

画面で学ぶ
ブラフボディーの流力振動

1. 円柱の流力振動

Flow-Induced Oscillations of Circular Cylinders

2. 矩形柱の流力振動

Flow-Induced Oscillation of Rectangular Cylinders

3. 振動流における物体の流体力学

Fluid Dynamics around a Bluff Body Submerged in Oscillatory Flow

4. 直列2円柱および2角柱の流力振動

*Flow-Induced Oscillation of Two Circular Cylinders and
Two Rectangular Cylinders in Tandem Arrangement*

— 流体力学の薦め —

ブラフボディーの流力振動 その4

直列2円柱と2角柱の流力振動

*Flow-Induced In-Line Oscillation of Two Circular Cylinders and
Two Rectangular Cylinders in Tandem Arrangement*

岡島 厚

Professor Emeritus, Kanazawa University

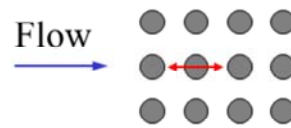
1. 矩形柱の静特性: Fluid-dynamic characteristics of rectangular cylinders

後流に渦が形成されるブラフ・ボディーのうち、種々な断面比の矩形柱の静特性、特にレイノルズ数変化と断面比に対する抗力係数、後流の渦ストローハル数の変化、そして矩形柱周りの流れパターンの変容について、実験結果やLESシミュレーションの結果を用いて説明する。

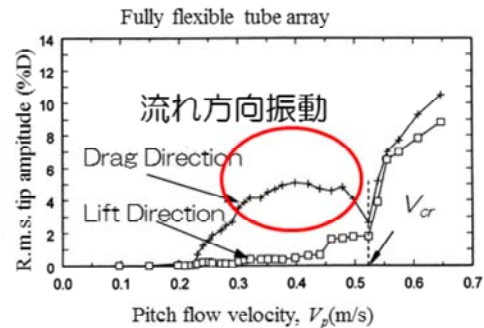
2. 矩形柱の流力振動: Flow-Induced Oscillation of rectangular cylinders

種々な断面比の矩形柱では、流れ方向振動、直角方向振動の渦励振やギャロッピングなどの自励振動(フラッター)が生じる。種々な矩形断面柱の流力振動現象を、実験結果やLESシミュレーション結果を用いて説明する。

各種プラントの配管内構造物などに利用されるブラフボディには、流力振動 (Flow-Induced Vibration) が発生する。



熱交換器など複数構造物が近接すると、単独物体と比べ流れ特性はより複雑になり、新たな流力振動が起こる。



直列に配置した2円柱の後方にスプリッタプレートを設置し、流れ方向振動について研究した。
風洞を用いて、一様流中に直列に配置した2円柱の後方にスプリッタプレートを設置し、下流側円柱を自由振動させ、振動模型の流れ方向振動特性と、流れの可視化から後流渦パターンとの関連を模型間隔を変えて、実験を行った。

◎複数物体周りの流れ現象の面白さ:

◎複数物体の静特性:

近年、熱交換器、送電線鉄塔、多脚煙突、橋梁、円筒状建築物および海洋構造物など、一様流中にある複数円筒状構造物の空力特性および空力弾性的不安定振動などが注目されている。本研究では、その最も基本的な場合として、流れに直列におかれた2本円柱の場合を取扱う。直列2本円柱まわりの流れに関する研究は、古く1933年Biermannら(1)による抵抗の測定以来、今日まで直列2本円柱の空力弾性的不安定性に関する研究なども含め、数多く成されている(2)–(12)。しかし、それらの実験の円柱直径に関するレイノルズ数は比較的 low、上流側円柱まわりの流れは亜臨界域である。しかし、実際にしばしば起こり得る、更に高いレイノルズ数範囲の直列2本円柱に関する実験的研究はほとんどない。一方、一様流中における単独円柱の場合でも、一般に臨界レイノルズ数以上の高いレイノルズ数範囲の風洞実験は難しい。そこで文献(13)では単独円柱表面に表面粗さを付けることにより、円柱表面の境界層の層流から乱流への遷移を促進させ、臨界レイノルズ数を低下させることを試みた。その結果、比較的低いレイノルズ数で円柱まわりの流れは超臨界域となり円柱表面に層流気ほうが形成される。その時の抵抗は極端に減少し、後流幅は狭くなり、後流の速度変動の卓越するストローハル数は亜臨界域に比べ大きくなる。そして更に層流気ほうが消滅し、円柱表面境界層が直接乱流はく離するような極超臨界域 (post-critical region) の流れが風洞実験で実現することができ、その時の円柱後流に極めて周期性の高いうず列が形成されることなどを明らかにした。すなわち

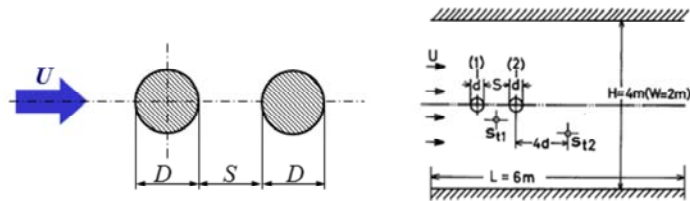
、高レイノルズ数における単独円柱まわりの流れおよび後流の様相は、臨界レイノルズ数以上でかなり変化する。従って、本研究の直列2本円柱の場合、上流側円柱まわりの流れが超臨界域や極超臨界域の時には、流れの様相は更に複雑になると考えられる。そこで本研究では、上流側円柱まわりの流れが超臨界域および極超臨界域であるような高いレイノルズ数において、直列2本円柱の抵抗や後流の速度変動の周波数などを測定し、その静的空力特性を明らかにするとともに、油膜法により円柱表面近傍の流れを可視観察し、流体力と流れの関連性を考察する。なお、文献(13)では単独円柱の場合の数種類の表面粗さにつき、その効果を調べたが、本実験ではそのうち最も代表的な2種類の表面粗さを用いた。

◎複数物体の流力特性：

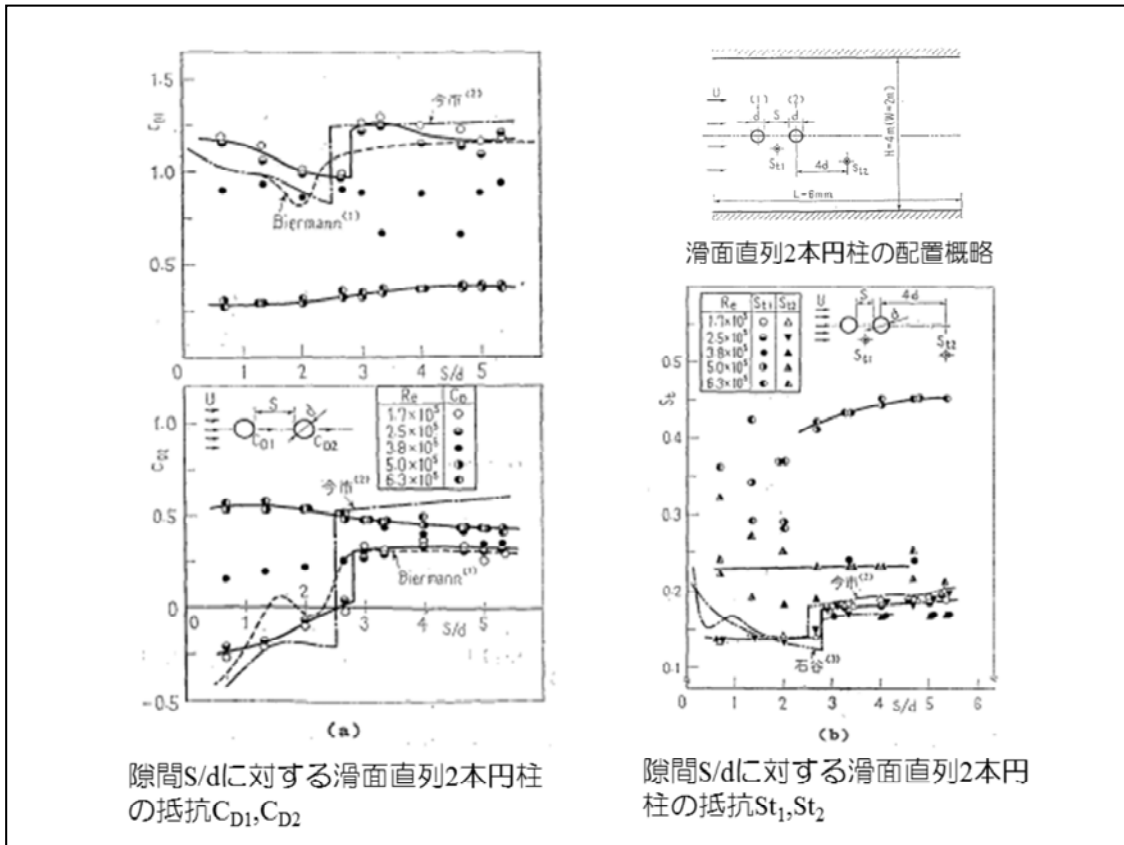
Two Tandem Cylinders in a Flow

(1) Statistic Characteristics of High Reynolds Numbers

(2) Aeroelastic Characteristics



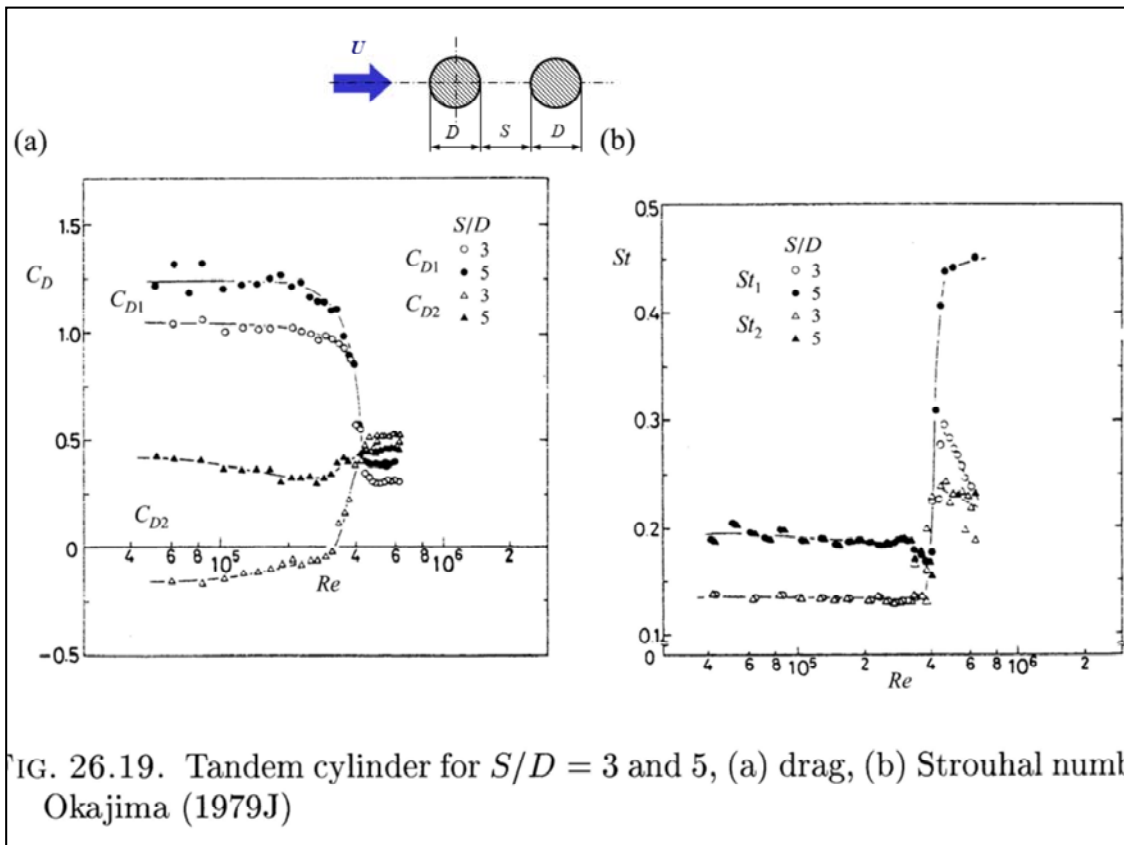
直列2本円柱

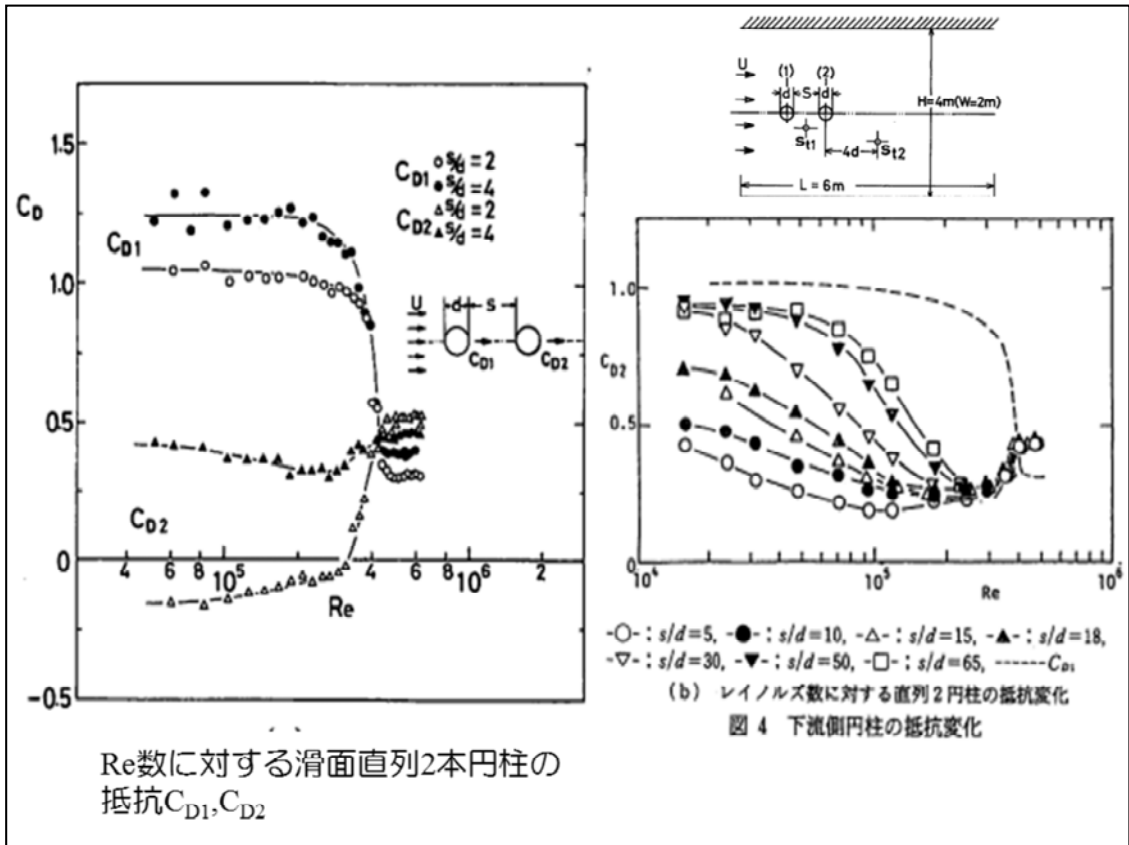


3. 1. 2間隔S/dの影響:

次に $Re=1.7 \times 10^5$, 2.5×10^5 (亜臨界), 3.8×10^5 (臨界域), 5.0×10^5 , 6.3×10^5 (超臨界域) の場合を例にして, 間隔 $S/d=0.7 \sim 5.3$ の範囲の (C_{D1} , C_{D2}) と (St_1, St_2) の変化を図5(a), (b) に示す. 図にはBiermannら(1)($Re=10^5$), 今市ら(2)($Re=2.1 \times 10^4$), 石谷ら(8)($Re=1.5 \times 10^3 \sim 10^4$) の値とも比較してある. これらの実験はいずれも亜臨界域のもので, 本実験の $Re=1.7 \times 10^5$, 2.5×10^5 の値とよい一致がみられる. 特に C_{D2} の本測定値は Re 数の最も近いBiermannらの値に一致している. また C_D や St 数の値が間隔 S/d に対し, ステップ状に変化する S/d の値は, この場合2.8である. 亜臨界域における C_D や St 数値のステップ状に変化する現象は, 上流側円柱背後のうず形成領域と下流側円柱の位極との相対的關係に密接に関連しており, これらに関して石谷ら(8)によるシュリーレン写真による可視観察やうずの周波数測定, 今市ら(2)による圧力の詳細な測定および渡辺ら(6)による流体力の測定などがある. これらの結果や本実験結果から亜臨界域で間隔の狭い場合, 両円柱間の速度変動はわずかで, 上流側円柱背圧と下流側円柱のよどみ点圧力はほぼ等しく, 下流側円柱背後のみうず列が生じ, 2本円柱流れはむしろ死水域で継がれた一つの物体まわりの流れの挙動を呈するといえる. そして $S/d=2.8$ を境にして流れは変化する. すなわち, 2本円柱間でも流れの周期的変動が強く感知されるようになり, C_D や St 数の値が急変する. 図5(b)に示すように, 亜臨界域の $S/d > 2.8$ の上・下流側円柱背後の St 数は, 共に $St_1 = St_2 = 0.19$ で単独円柱の場合に等しい. また C_D は, 図5(a)から, $C_{D1} = 1.2$ で単独円柱の値に近く, $C_{D2} = 0.3$ でかなり小さい. 次に, 臨界域($Re=3.8 \times 10^5$)の場合, C_D や St 数の値は $S/d=0.7 \sim 5.3$ にわたり, 亜臨界域と次に述べる超臨界域における値の中間的傾向を示し, S/d の影響はわずかで, 特に St 数の値のばらつきが目立つ. そして, 上流側円柱に層流気ほうの形成されている超臨界域

($Re=5.0 \times 10^5, 6.3 \times 10^5$)の場合,図5(a), (b)に示すように, 両円柱の(C_{D1}, C_{D2})と(St_1, St_2)の値に及ぼす S/d の影響は少なく, わずかに S/d の増加に伴い, C_{D1} は漸増するのに対し, C_{D2} は漸減し, 両者の値が接近しているが, 本実験範囲では, C_{D1} より C_{D2} のほうが常に大きな値である. この場合, 上流側円柱における層流気ほうの形成により C_{D1} の値は亜臨界域に比べ減少する. 同時に上流側円柱背後の逆流領域が極端に縮小し, 後流幅が狭くなることにより, CD 図は逆に増大する. しかも図3(b)に示したごとく, 下流側円柱の流れには層流気ほうは無く, 約 $\theta=120^\circ$. 付近で乱流はく離している. 図6には油膜法によって求めた超臨界域($Re=5.0 \times 10^5$)における流れのはく離点や, 再付着点の位歴の S/d に対する変化を示す. 図から上流側円柱の層流気ほうの位歴や上・下流側円柱のはく離点の位図などは $S/d=0.07 \sim 4.67$ の範囲であり変わらない. このことは, 図5(a)の超臨界域において C_{D1} や C_{D2} の値が S/d に対し, 大きな変化がないことに対応している.





上・下流側円柱に作用する抵抗係数 $C_{D1(2)} = 1/2 \rho U^2 J(D)$ (D は抵抗, ρ は空気密度, 添字1は上流側, 2は下流側を示す), および上・下流側円柱背後のストローハル数(s' 数) $S''_{(2)} = \text{Gal}(2)/U(\alpha, (2))$ (α は速度÷変動の周波数), そして2本円柱まわりの流れにつき, レイノルズ数(R ・数)や2本円柱の間隔 S/α を変え, それらの影響を議論する.

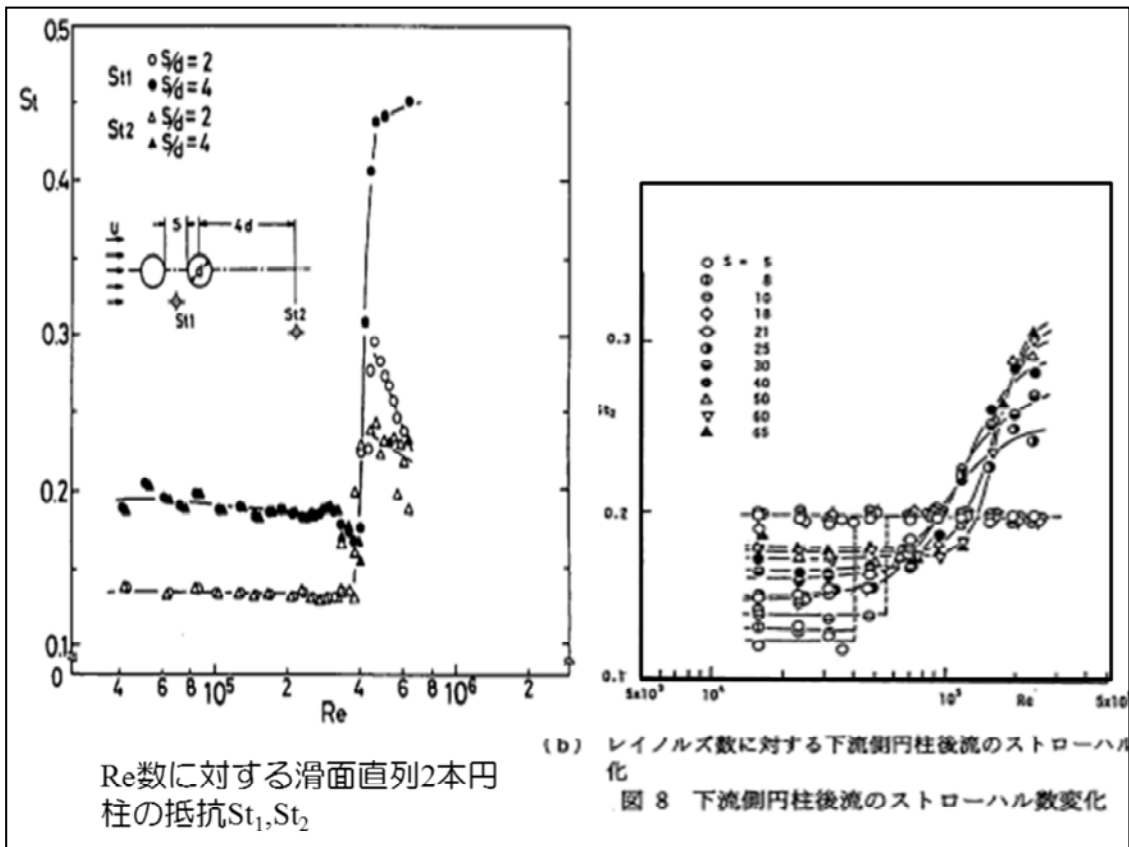
3. 1滑面円柱の場合

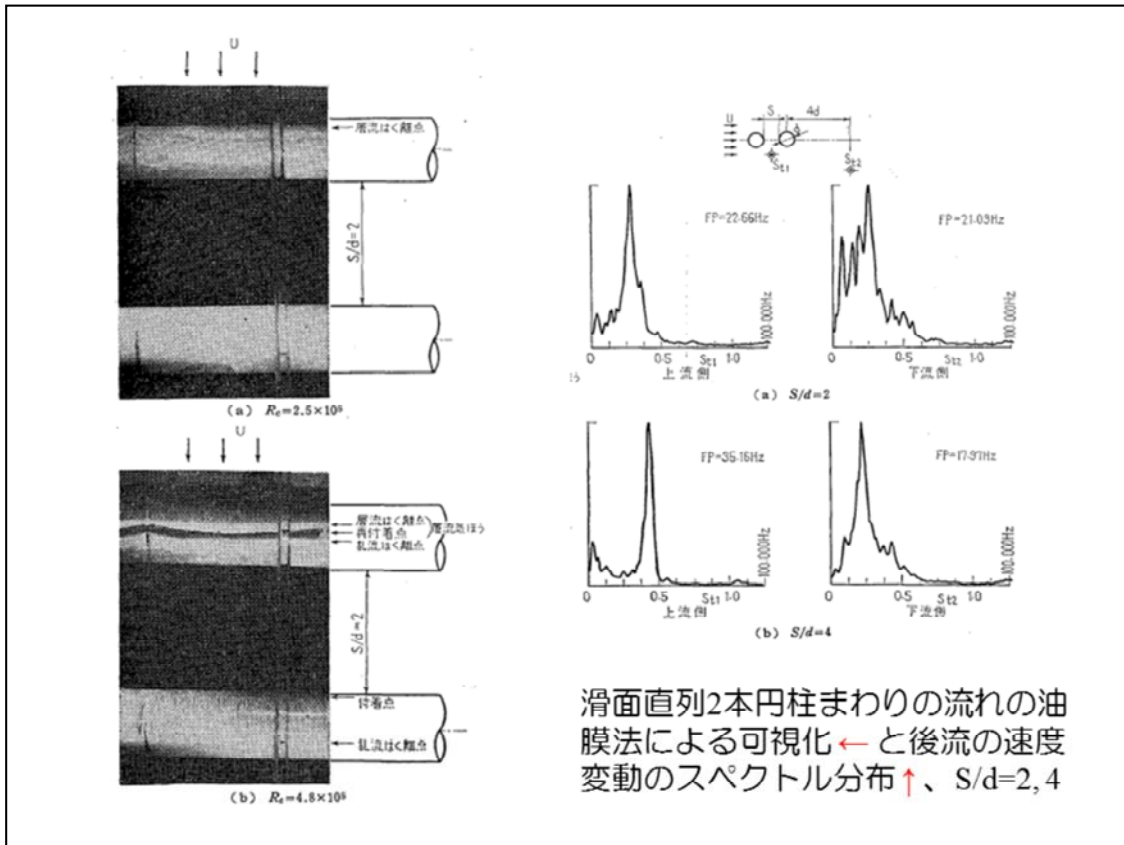
3.1.1 R 数の影響円柱間隔 $S/d=2, 4$ の場合

を例にして, $R=0.53 \times 10^5 \sim 6.3 \times 10^5$ の範囲の直列2本滑面円柱に働く抵抗(C_{D1}, c_{D2})および S_1 数(S'' 函)の変化を図2(a), (b)に示す. 図によれば, 上流側円柱の C_{D1}, S'' の値は R ・数が約 3×10^5 に達するまで大きな変化はなく, $C_{D1} = 1.0(S/\alpha - 2), 1.2(S/\alpha - 2)$ ($d=4$), $S'' = 0.14(S/\alpha - 2)$, 速度変動の振幅は非常に小さい), $0.19(S/\alpha - 4)$ である. そして臨界 R 数の $R_c = 3.8 \times 10^5$ 付近で c_{D1}, S'' の値は急激に変化し, $R > 4.4 \times 10^5$ の範囲では $C_{D1} = 0.3(S/\alpha - 2), 0.4(S/\alpha - 4)$ でかなり小さい値となる. 一方この時の卓越す

るS`数の値は大きくなる. 結局, $S/Cf=2, 4$ の場合共に, 上流側円柱の抵抗やS`数のR`数に対する変化は単独円柱の場合(8)とほぼ同様の様相を呈し, その臨界R`数の値も $R_{crit}=3.8 \times 10^5$ で単独円柱と変わらない.

一方, 下流側円柱の CD_2, S_2 の値は主として上述のような上流側円柱まわりの流れの変化に付随して変わる. すなわち, $R < 3.0 \times 10^5$ のR`数範囲では CD_2 は推力 $CD=-0.1 \sim 0$ ($S/\alpha=2$)あるいは比較的小さい値 $CD_2=0.4 \sim 0.3$ ($S/\alpha=4$)で, S`動の値は上流側の S_n と同じである. そして CD_1, S_1 の値が急変する臨



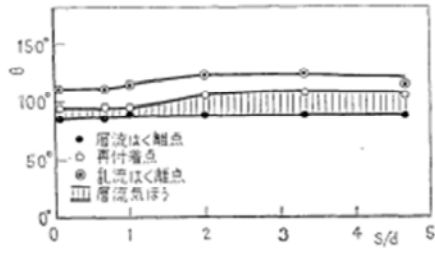


滑面直列2本円柱まわりの流れの油膜法による可視化 ← と後流の速度変動のスペクトル分布 ↑、 $S/d=2, 4$

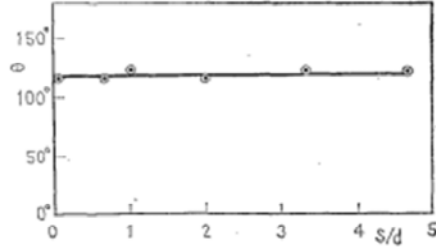
界 R 僅数付近で CD_2 , s :量の値も同時に変化し, $R' >$
 ,好 1×10 , では, $111]$ 開 S/d にほとんど依存せず, $cD_2 =$
 0.5 で, 上流側円柱の cD_1 よりむしろ大きな値であ
 る.

次に, 円柱近換流れの抽膜法による可視観察結果を
 図3(a),(b)に示す. 図3(a)($R = 25 \times 10^3$, S/d
 $= 2$)の上流側円柱においては, 白色の油膜が吹含取
 られ, 黒色の円柱表面が露出している部分と抽膜の残
 っている部分に減別され, その塊界はよどみ点から約
 $\theta = 70^\circ$ の位肚で, この付近から塊界層は圃流はく離
 している. 上流側円柱まわりの流れは亜臨界域であ
 る. 下流側円柱の油膜模様は判然としない. 一方, 図
 3(b)($R = 1.8 \times 10^3$, $S/d = 2$)の場合, 上流側円柱
 の油膜模様は四つの部分に分かれている.よど象点か
 ら $\beta = 90^\circ$ までの油談の筈い部分, $\theta = 90^\circ \sim 105^\circ$ の
 白い抽膜の残っている部分, $\beta = 105^\circ \sim 120^\circ$ の油膜の
 全く無い部分, そして $> 120^\circ$ の抽膜の残っている
 部分である. しかも, この油膜換槻はスパン方向にほ
 ぼ直線で, 流れは二次元的である.この場合, 流れは

$\beta=90^\circ$ で層流ばく雄するが、 $\beta=105^\circ$ で再付着して、層流気ほうを形成し、 $\beta=120^\circ$ で乱流はく離している。流れは超臨界域である。下流側円柱の流れは 0.30 。付近から円柱に沿い、 0.125 で乱流ばく離していることが観察される。図4(a)にはこの時の円柱後流のパワースペクトル密度分布を示すが、上・下流側の S' 数、 $S_1 \dots S_{10}$ の帯域はかなり広く、その卓越する S' 数はそれぞれ梱逸している。さらに、図4(b)には $S'_i - 1$ の掛合を示すが、上流 0.45 の卓越する S' 数は、 $S'_3 = 0.45$ で11i独円柱の超臨界域の場合に一致している。そしてこの周波数成分は比較的早く減衰し、下流側円柱ではその成分はわずかに検出される極度で、下流側の S'_i 数は S'_1 と異なり、 $S'_1 = 0.23$ である。以上の結果から直列2本円柱まわりの流れの R' 数に対する変化は、主として上流側円柱流れの変化に支配されていると言える。すなわち臨界 R' 数以上では、単独円柱同様上流側円柱に層流気ほうが形成され、乱流はく離点は後退し、上流側円柱背後の逆流領域の大きさが縮小し、しかも後流幅も狭くなる。このような流れの変化に対応して下流側円柱まわりの流れは変化する。従って図2(a),(b)に示すごとく、直列2本円柱の (CD_h/CD_2) や (S'_1/S'_2) の値は上流側円柱の臨界 R' 数 0.45 付近で同時に変化している。そこで以後、直列2本円柱の上流側円柱まわりの流れが亜臨界域、臨界域、超臨界域、極超臨界域であることを、便宜上単に亜臨界域、臨界域、超臨界域、極超臨界域と呼ぶことにする。

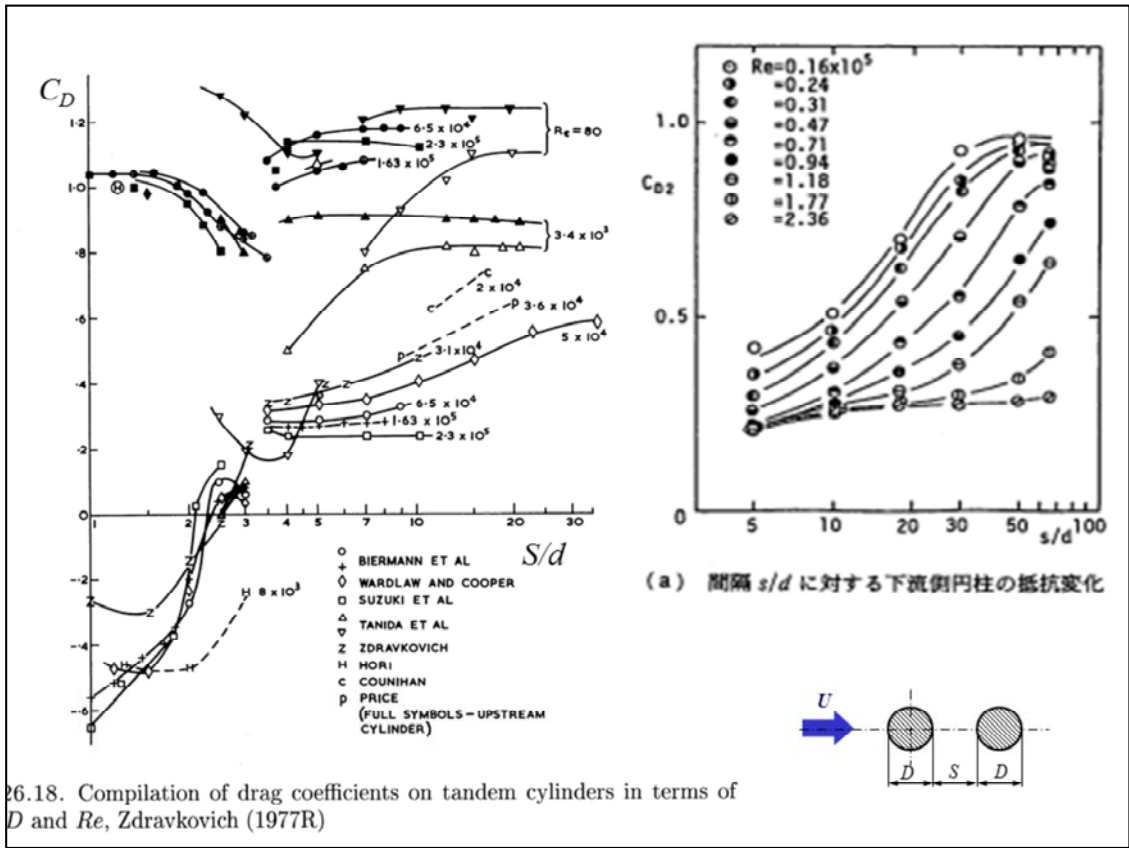


(a) 上流側円柱



(b) 下流側円柱

滑面直列2本円柱の剥離点および再付着点の位置、 $Re=5.0 \times 10^5$



26.18. Compilation of drag coefficients on tandem cylinders in terms of D and Re , Zdravkovich (1977R)

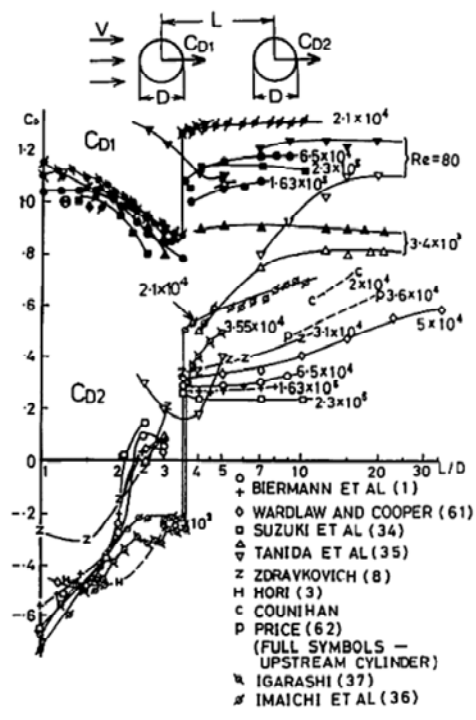
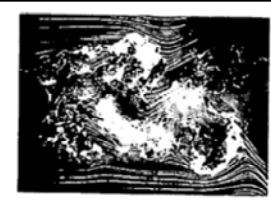
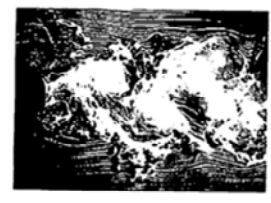


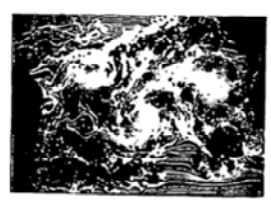
Figure 29. Drag coefficients of two tandem circular cylinders [7, 35-37].



(A) $L/D = 6$

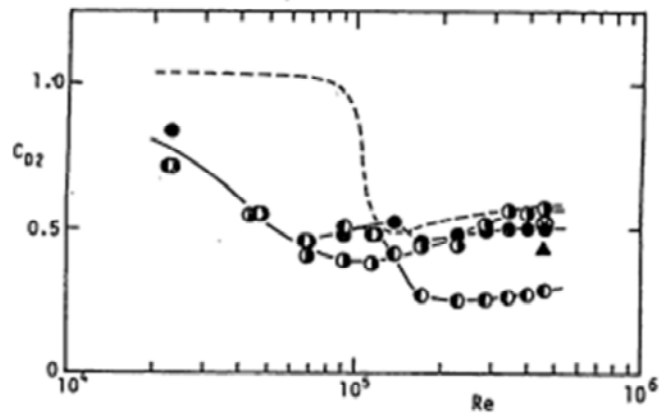


(B) $L/D = 11$



(C) $L/D = 21$





$s/d=20$ { \circ : 滑面-あらし, \circ : あらし-滑面
 \bullet : あらし-あらし
 $s/d=10$, \blacktriangle : あらし-あらし
 $s/d=29$, \circ : あらし-あらし
 : 上流側表面あらし円柱の C_{D1}

図 10 直列表面あらし円柱の抵抗

Roughed Circular Cylinders

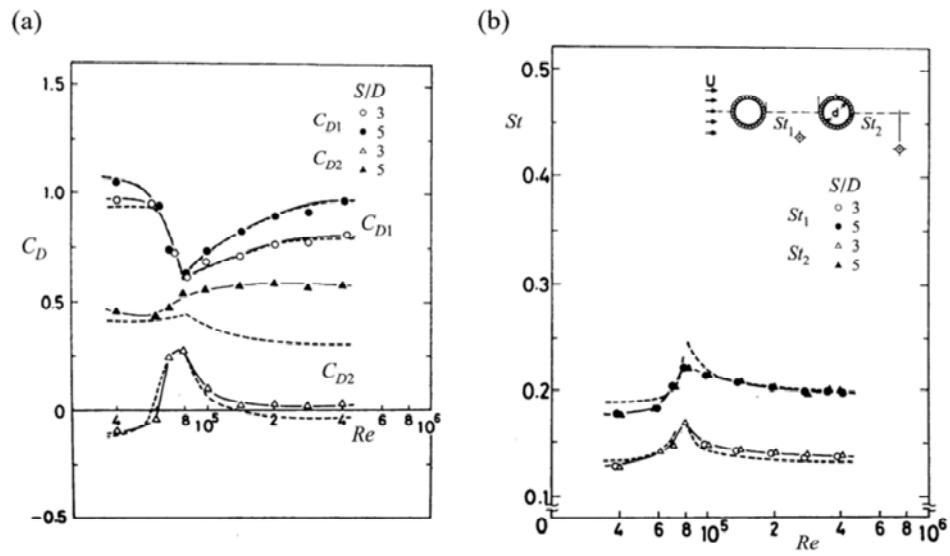
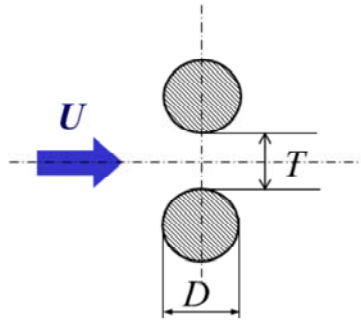


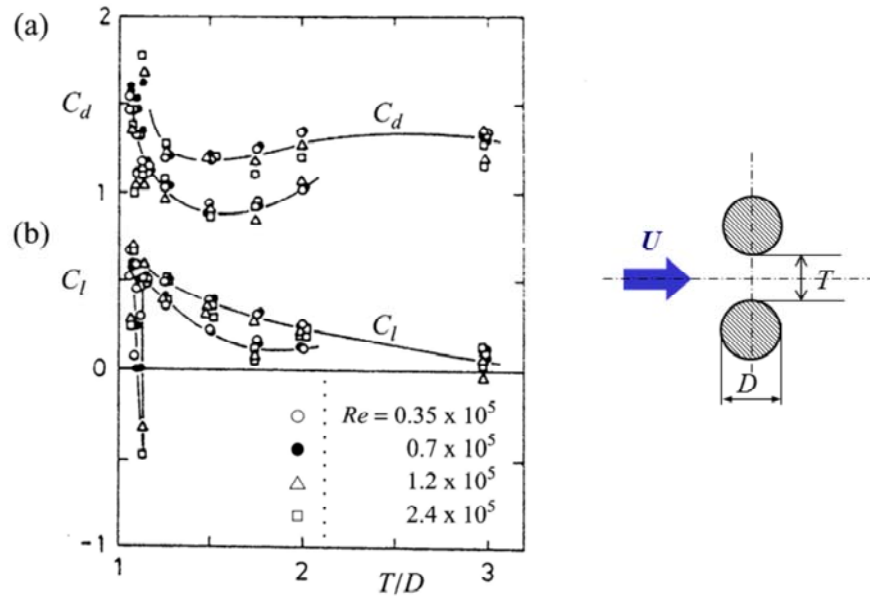
FIG. 26.24. Effect of surface roughness $K/D = 0.9\%$ in terms of Re , on (a) drag, (b) Strouhal number, Okajima (1979J)

Side by Side Cylinders in a Flow

(1) Statistic Characteristics

(2) Aeroelastic Characteristics





45. (a) Local drag, (b) local lift in terms of T/D for side-by-side cylinders at $Re = 120k$ and $240k$, Okajima *et al.* (1986J)

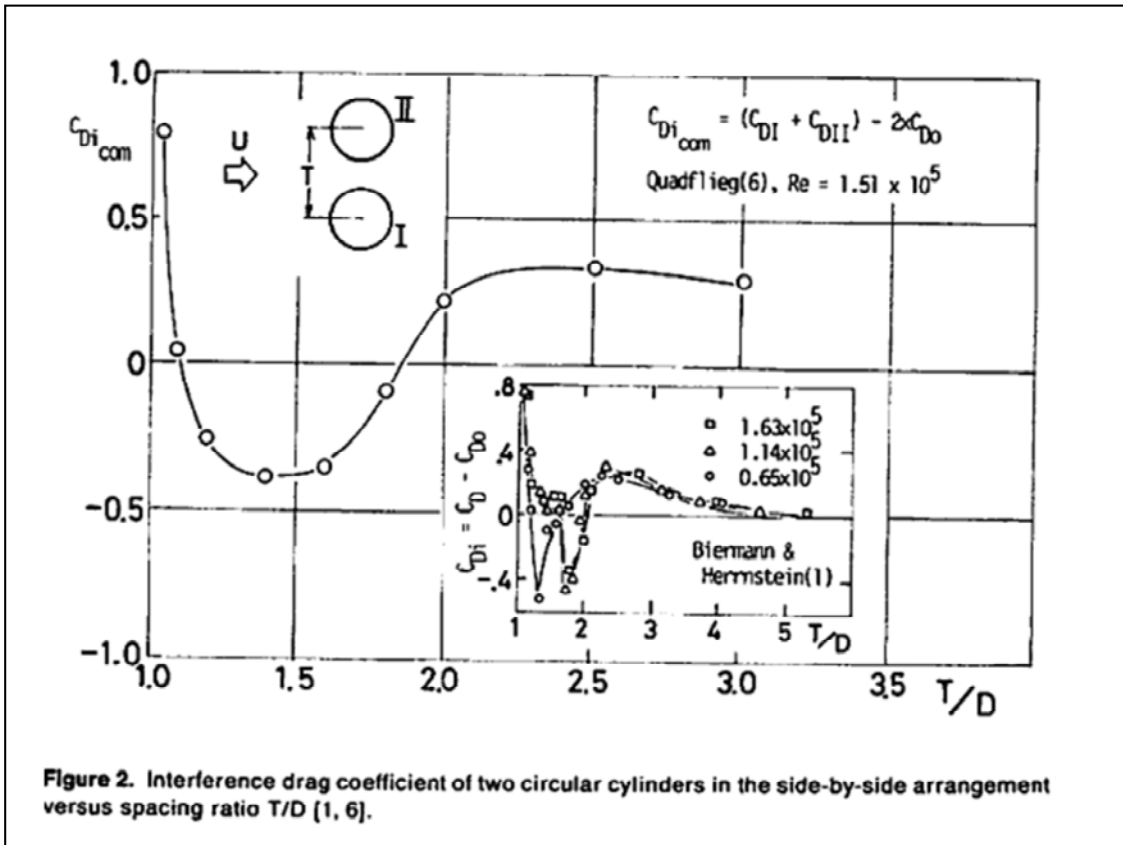


Figure 2. Interference drag coefficient of two circular cylinders in the side-by-side arrangement versus spacing ratio T/D [1, 6].

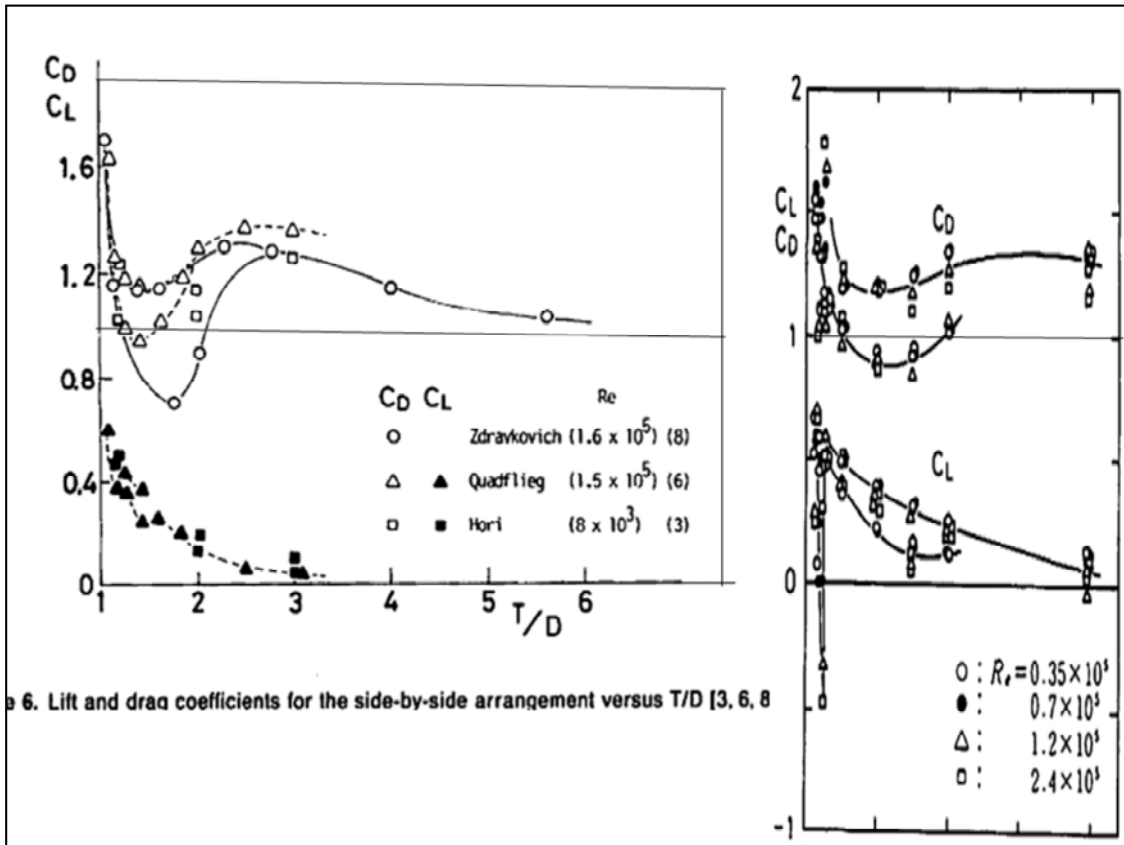
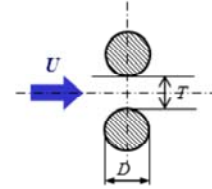
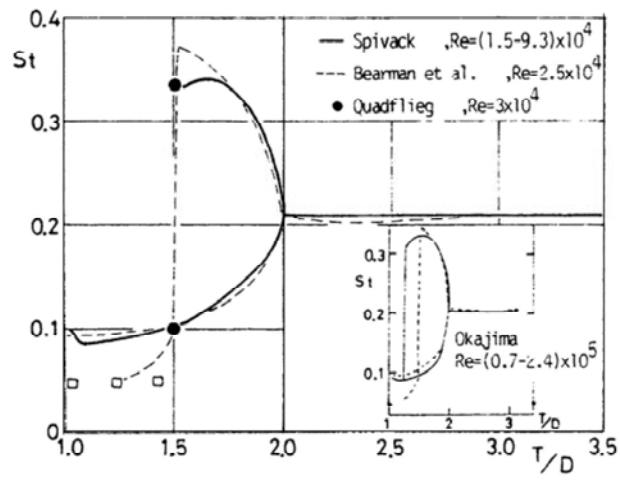
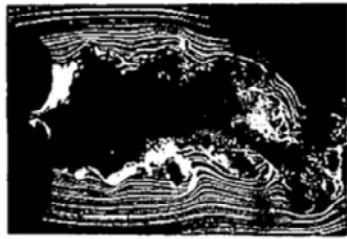


Fig. 6. Lift and drag coefficients for the side-by-side arrangement versus T/D [3, 6, 8]



G. 26.35. Strouhal number in terms of T/D , Okajima *et al.* (1986J)



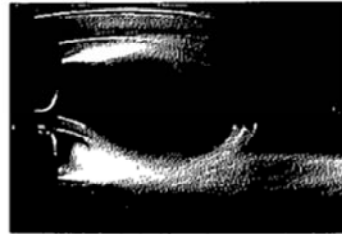
(A) $T/D = 1.125$



(C) $T/D = 3.0$



(B) $T/D = 1.5$



(B') $T/D = 1.5$
(Exposure Time=1sec)

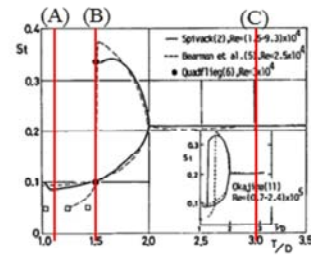


Figure 3. Strouhal number for the side-by-side arrangement versus T/D [2, 5, 6]

Figure 4. Flow patterns around two circular cylinders in the side-by-side arrangement at various spacing ratios [11, 23].

- (A) $T/D = 1.125$ (instantaneous flow).
- (B) $T/D = 1.5$ (instantaneous flow).
- (B') $T/D = 1.5$ (time-mean flow).
- (C) $T/D = 3$ (instantaneous flow).

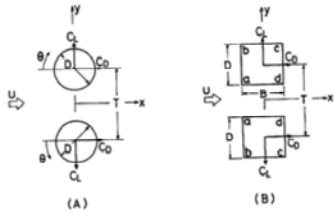


Figure 1. Side-by-side arrangement and definition of fluid force coefficients: (A) circular cylinders, (B) rectangular cylinders.

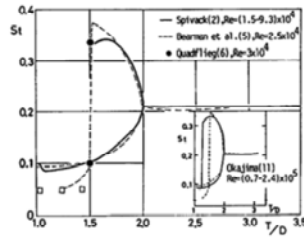


Figure 3. Strouhal number for the side-by-side

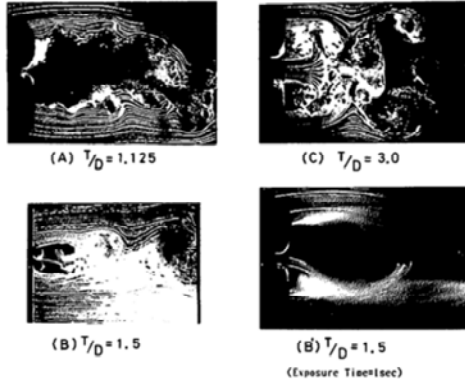


Figure 4. Flow patterns around two circular cylinders in the side-by-side arrangement at various spacing ratios [11, 23].
 (A) $T/D = 1.125$ (instantaneous flow).
 (B) $T/D = 1.5$ (instantaneous flow).
 (B') $T/D = 1.5$ (time-mean flow).
 (C) $T/D = 3$ (instantaneous flow).

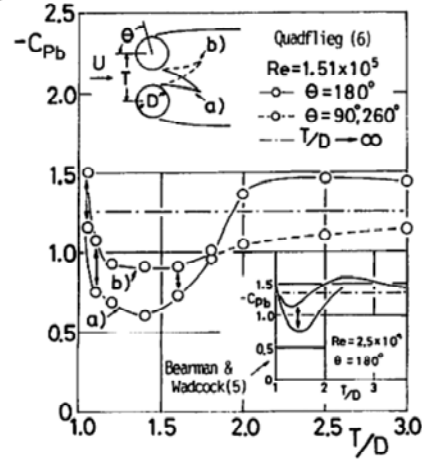


Figure 5. Coefficient of pressure vs T/D .
 Legend:
 - \circ Quadflieg (6), $Re=1.51 \times 10^5$, $\theta=180^\circ$
 - \dots Quadflieg (6), $Re=1.51 \times 10^5$, $\theta=90^\circ, 260^\circ$
 - $---$ Quadflieg (6), $Re=1.51 \times 10^5$, $T/D \rightarrow \infty$
 - \circ Bearman & Wadcock (5), $Re=2.5 \times 10^4$, $\theta=180^\circ$

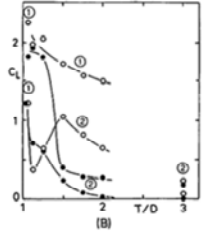
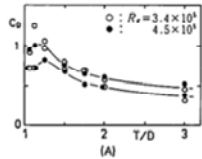


Figure 8. Lift and drag coefficients for the side-by-side arrangement of spherical Reynolds numbers versus T/D [11].
 (A) Drag coefficient.
 (B) Lift coefficient. The number shows the number of separation bubbles for each cylinder.

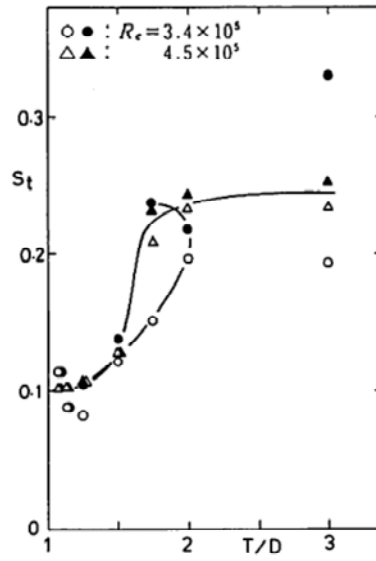
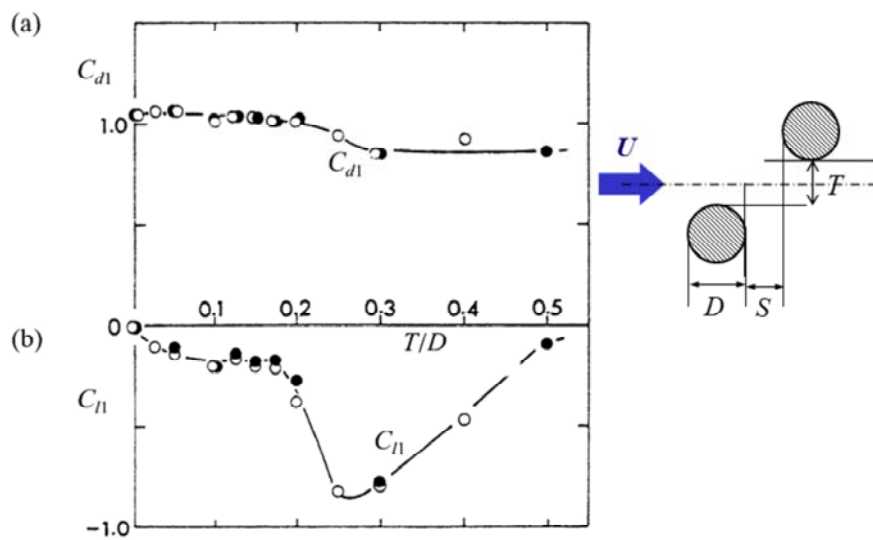


Figure 9. Strouhal number for side-by-side arrangement of spherical Reynolds numbers [11].



6.54. Variation in local (a) drag, (b) lift in terms of T/D for $S/D = 1.25$
 $Re = 100k$, Okajima and Sugitani (1976P)



Flow-Induced In-Line Oscillation of Two Circular Cylinders in Tandem Arrangement

Atsushi OKAJIMA,
Professor, Kanazawa-Gakuin College
and
Satoru YASUI, Takahiro KIWATA, and Shigeo KIMURA

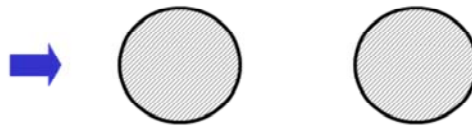
***Keywords:* flow-induced oscillation, in-line oscillation, two circular cylinders, tandem arrangement, flow visualization**

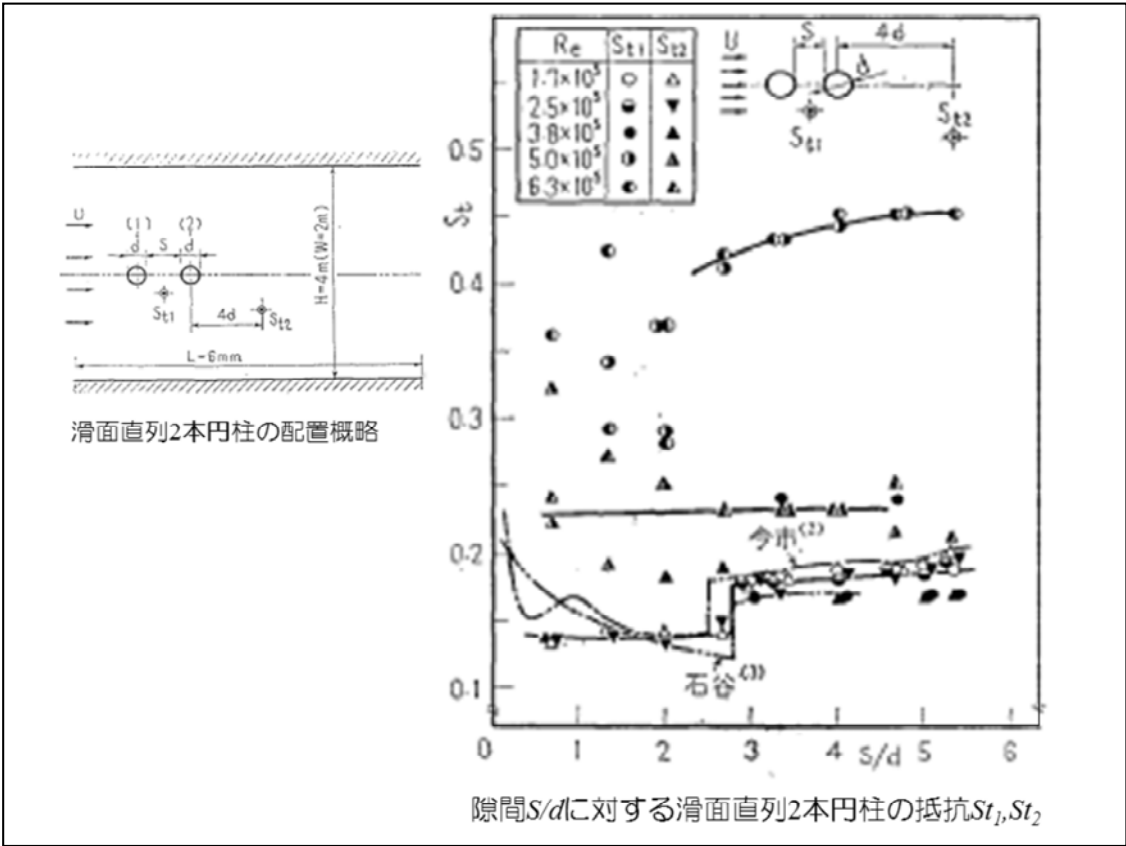
I would like to talk about Flow-Induced In-Line Oscillation of Two Circular Cylinders in Tandem Arrangement.

Two Tandem Cylinders

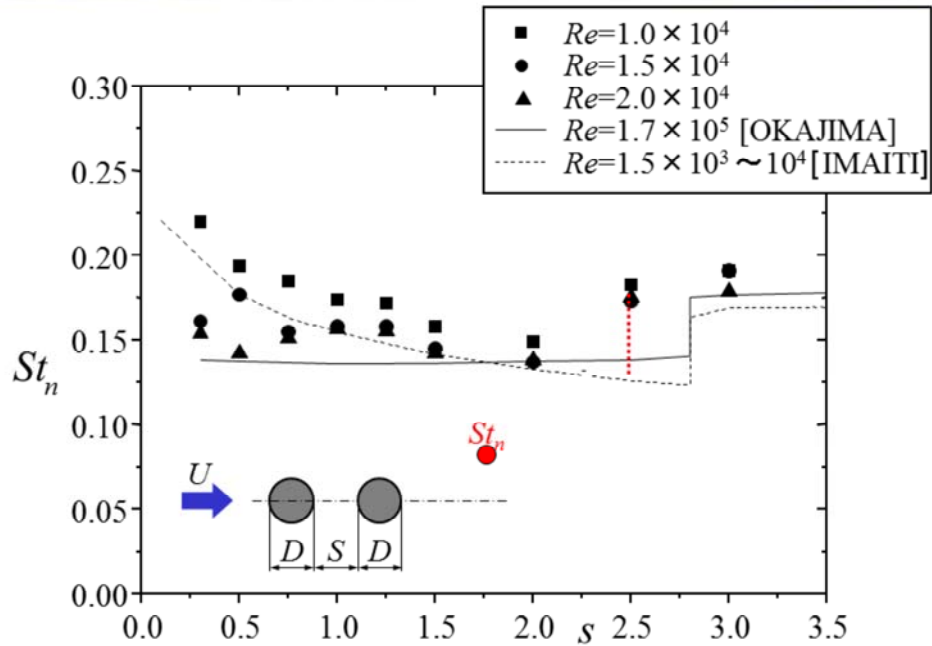
(1) Statistic Characteristics of High Reynolds Numbers

(2) Aeroelastic Characteristics





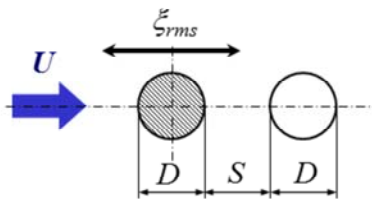
FUNDAMENTAL DATA : Strouhal Numbers of Two Stationary Cylinders in Tandem Arrangement



Strouhal number changes with spreading the gap-distance and is influenced by Reynolds number

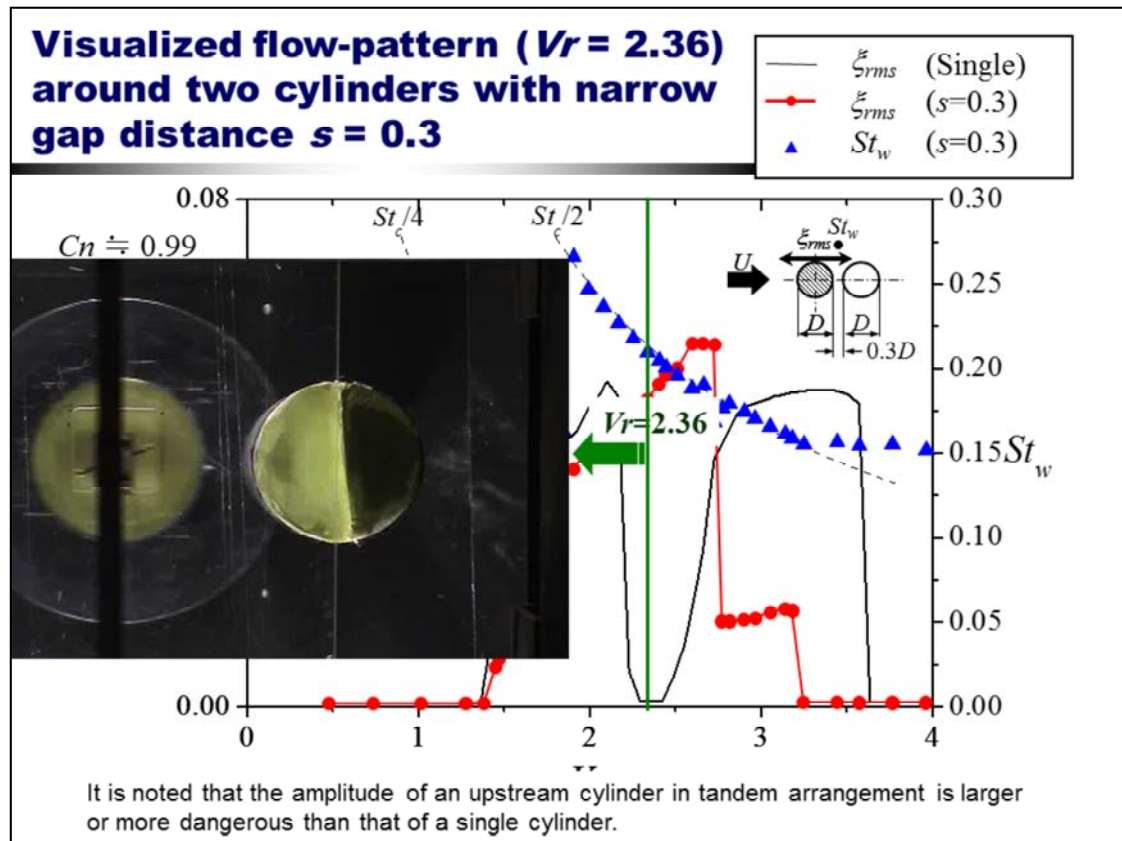
(6) And also we show the results of the Strouhal number behind two stationary cylinders in tandem arrangement against the gap-distance **as the fundamental data of two tandem cylinders**. You can see that the Strouhal number changes with spreading the gap-distance and is influenced by Reynolds number, also.

RESULTS: Flow-Induced In-Line Oscillation of the **Upstream** Cylinder

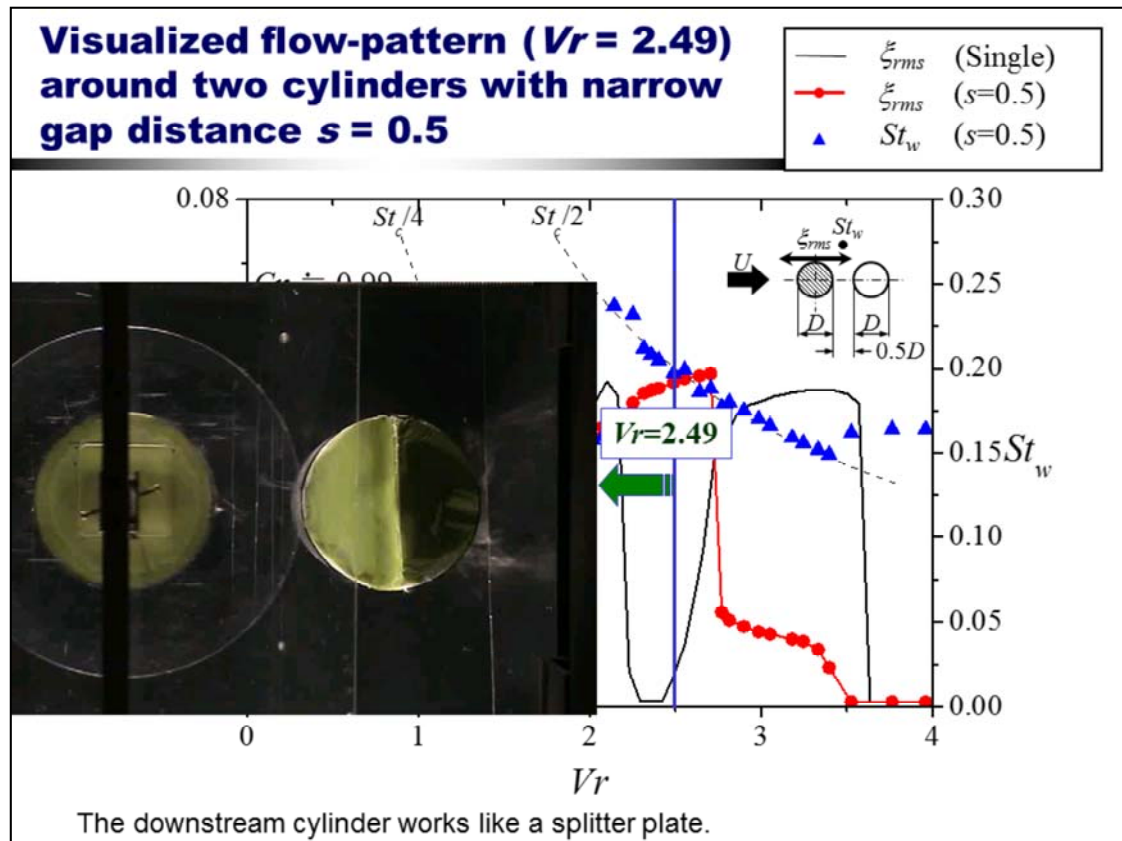


The upstream cylinder is free for in-line oscillation and the downstream cylinder is fixed.

(8) Well we present experimental results when the upstream cylinder is free for in-line oscillation and the downstream cylinder is fixed.

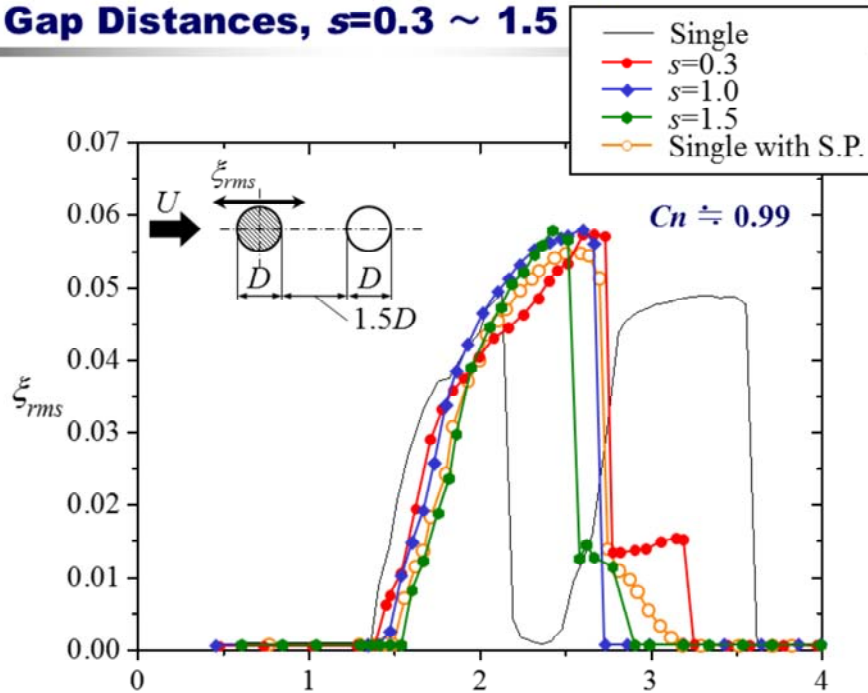


(9) This figure shows an example of the response curve of the in-line oscillation of the upstream cylinder for the narrow gap distance of 0.3. This figure shows response amplitude and the Strouhal number of a wake against reduced velocity, compared with the results for a single cylinder. You will find only one excitation region, which is different from two excitation regions for a single cylinder. We present a visualized animation for the upstream cylinder oscillating at reduced velocity 2.36, You can see, the shear layer separating from the upstream cylinder clearly rolls up into symmetrical vortices in the gap-field. It is evident that this in-line oscillation is induced by the symmetric vortices formed in the gap-field.



(10) This figure also shows another example of the response curve and the Strouhal number of the upstream cylinder oscillating for the narrow gap distance of 0.5. You can see a single excitation region like this. And we present a visualized animation for the upstream cylinder oscillating at reduced velocity 2.49. It is evident that this in-line oscillation is induced by the symmetric vortices formed in the gap-field.

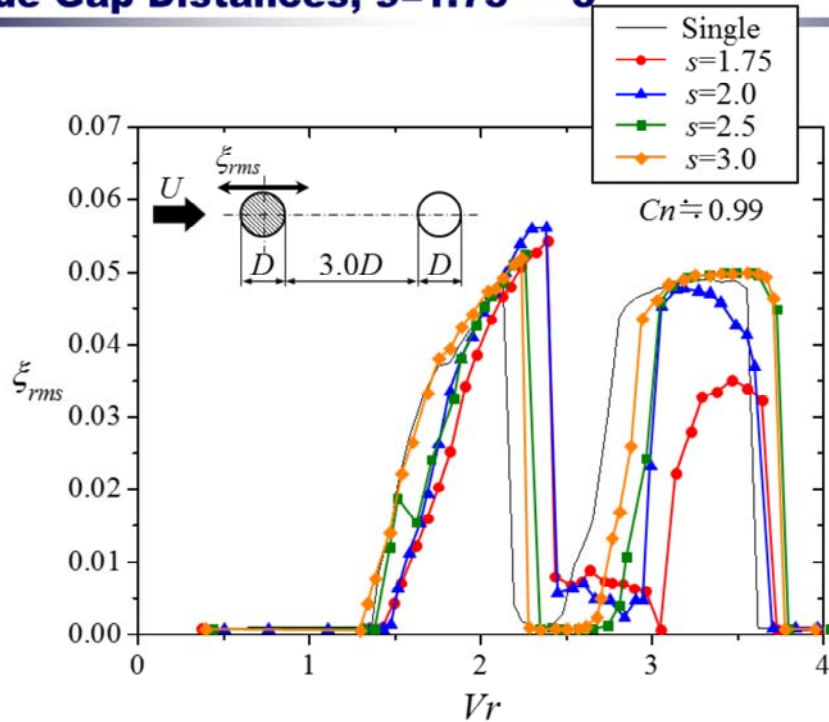
In-Line Oscillation of the Upstream Cylinder for Narrow Gap Distances, $s=0.3 \sim 1.5$



And it is noted that the maximum amplitudes of a cylinder with a splitter plate and two tandem cylinders, all are larger than that of a single cylinder.

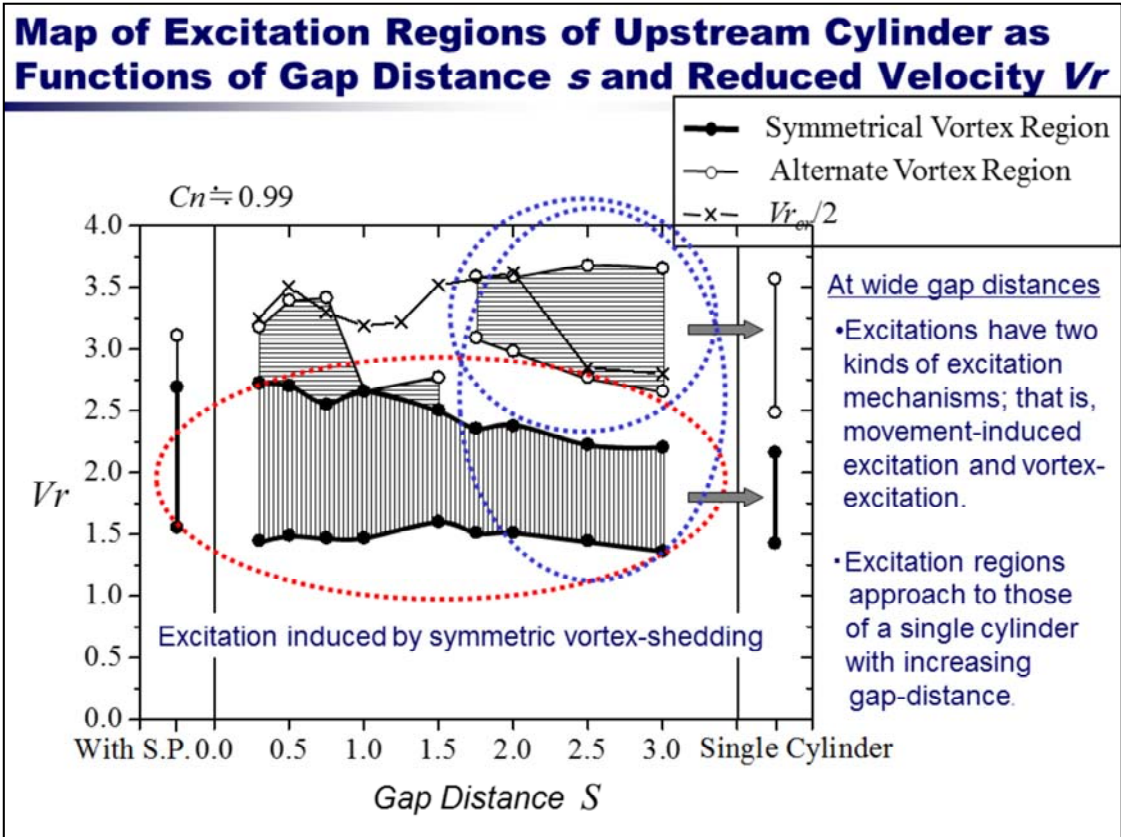
(11) This figure summarizes the response characteristics of the in-line oscillation of the upstream cylinder for the narrow gaps from 0.3 to 1.5. This is the response curve of a single cylinder. This is for a cylinder having a splitter plate. The gap distance changes into 0.3, 0.5, 1.0 and 1.5. You can see all response curves are similar to the response-characteristics of a cylinder with a splitter plate. It means that downstream cylinder works as a splitter plate.

In-Line Oscillation of the Upstream Cylinder for Wide Gap Distances, $s=1.75 \sim 3$



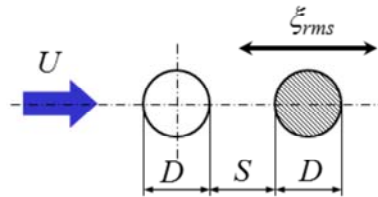
(12) The response curves for the wide gap distances from 1.75 to 3 are presented here.

It is evident that all the response curves gradually approach that of a single cylinder with spreading gap distance over 1.75.



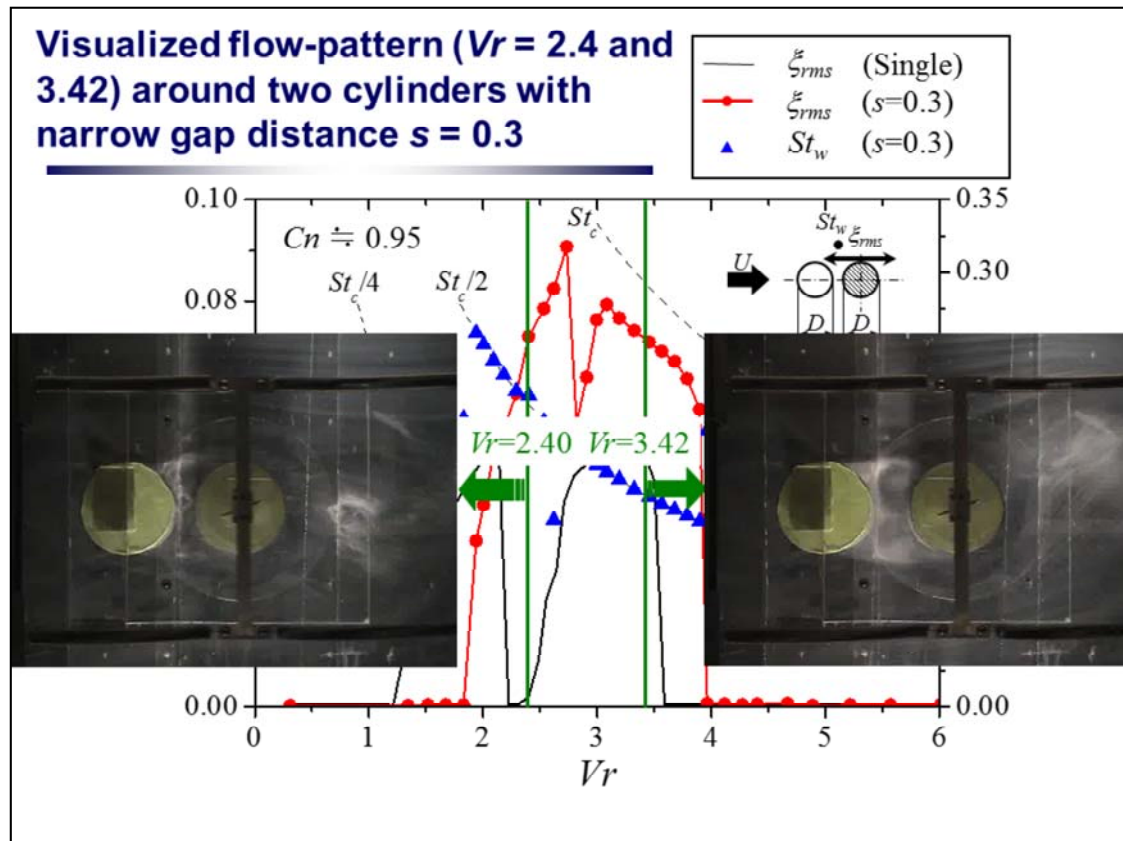
(13) Then we try to make a map of excitation regions of the upstream cylinder as functions of gap distance and reduced velocity. We can see two kinds of excitation regions due to different oscillation mechanisms; that is, symmetric vortices and alternate vortices. Excitation induced by symmetric vortex-shedding appears at all experimental gaps. The excitation of the upstream cylinder with wide gap have two kinds of excitation mechanisms; that is, movement-induced excitation and vortex-excitation. These characteristics approach to that of the single cylinder with increasing gap-distance over 1.75.

RESULTS : Flow-Induced In-Line Oscillation of the **Downstream** Cylinder

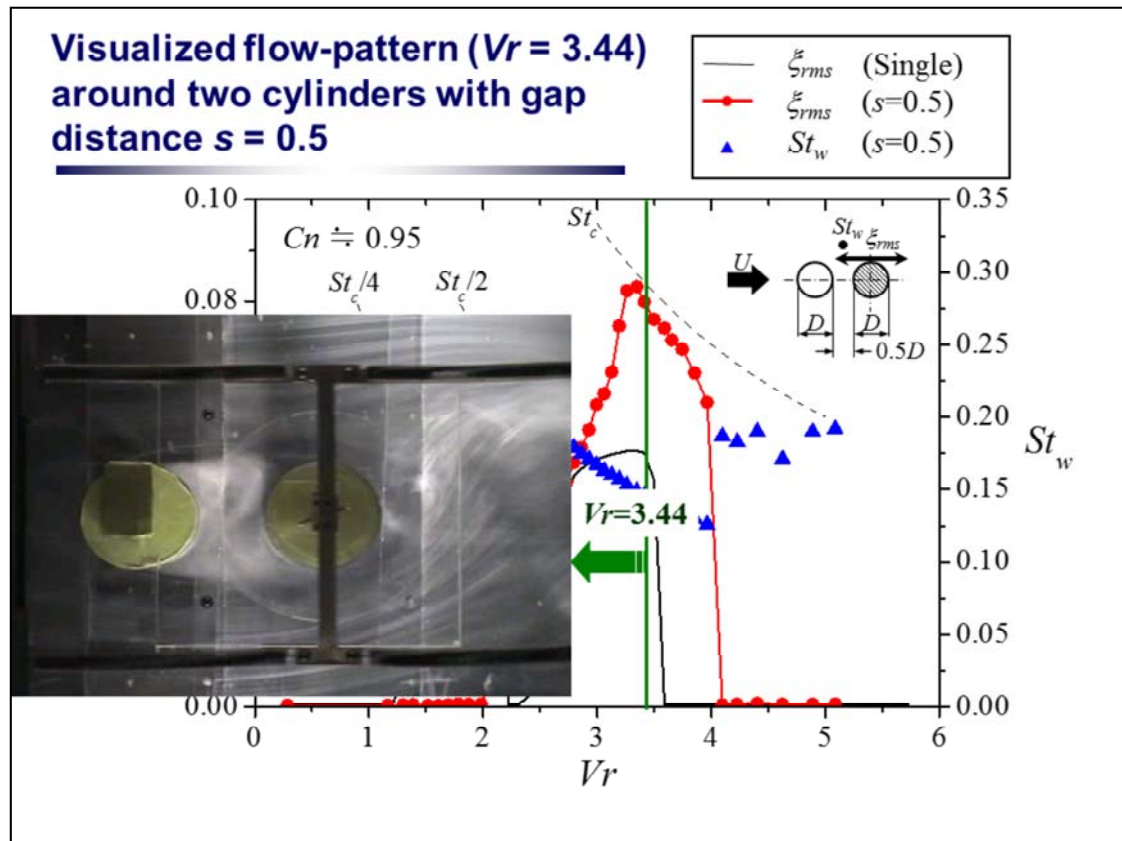


The downstream cylinder is free for in-line oscillation and the upstream cylinder is fixed.

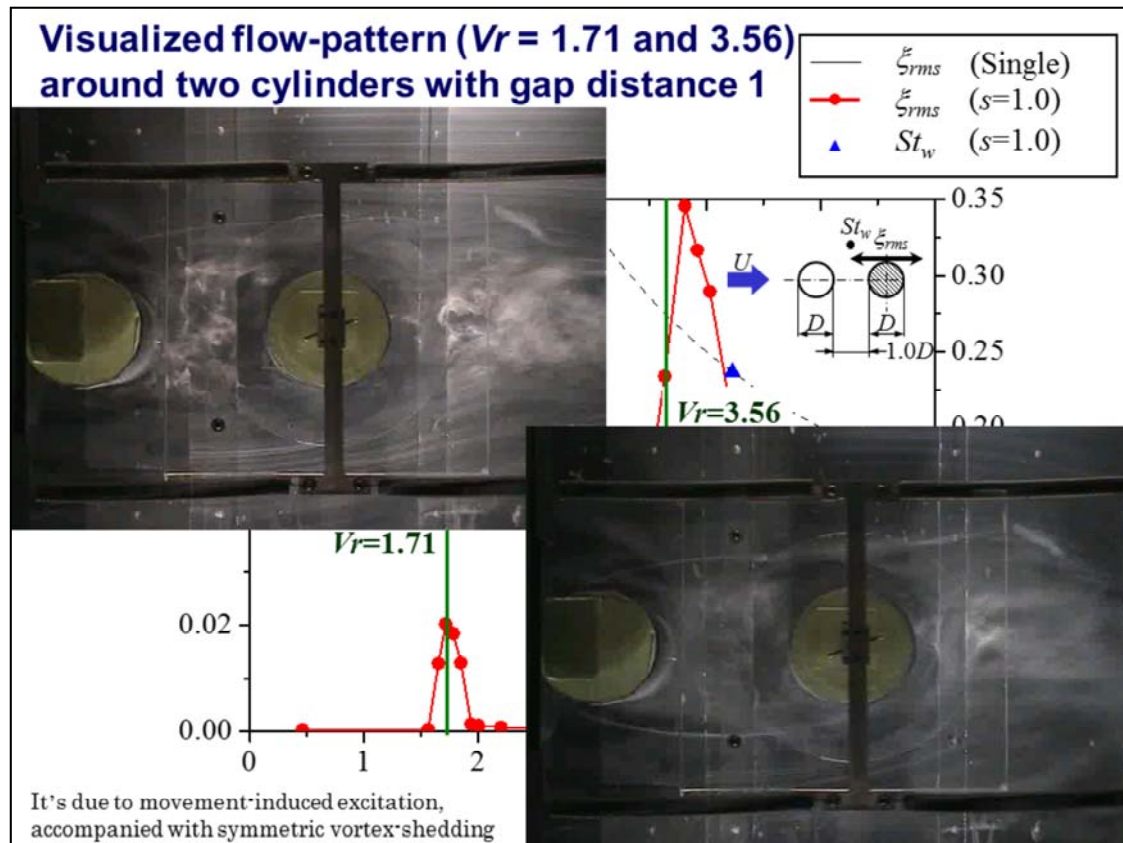
(14) Next we present the results when the downstream cylinder is free for in-line oscillation and the upstream cylinder is fixed.



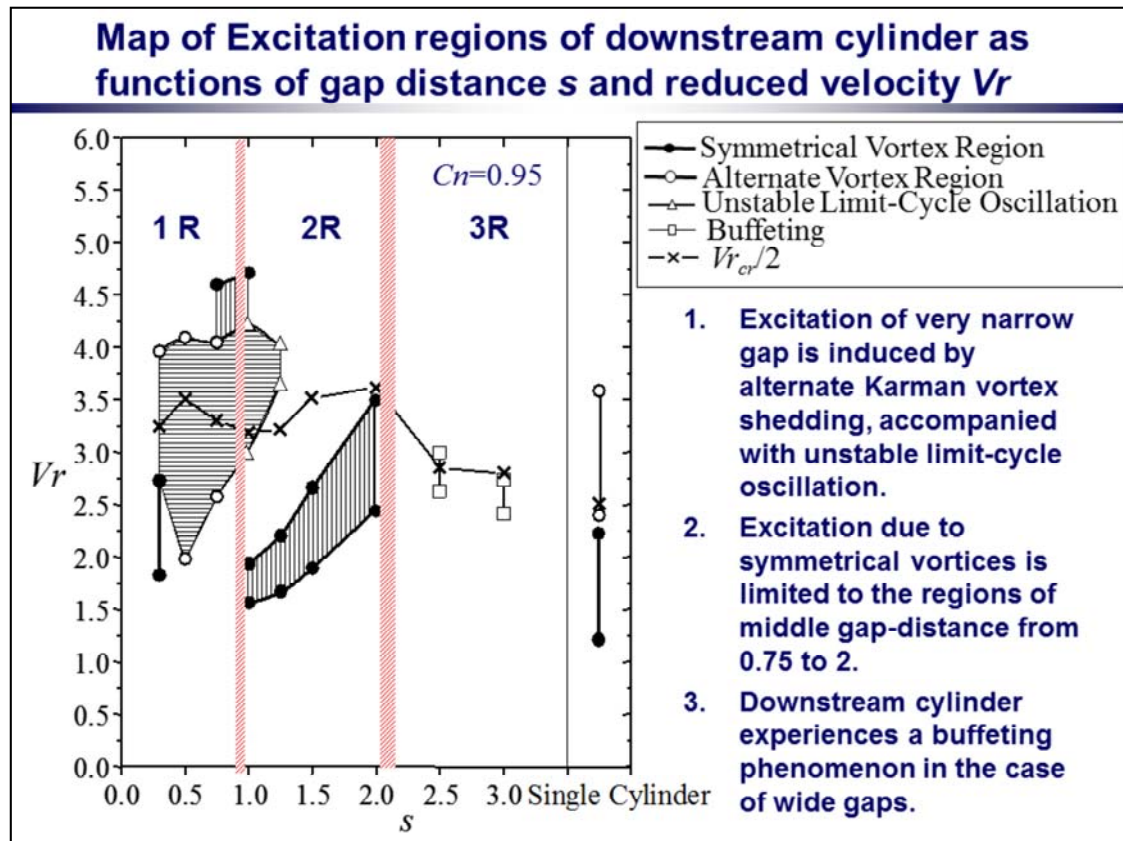
(15) This figure shows the response curve and Strouhal numbers of the downstream cylinder with narrow gap distance ratio of 0.3. The response characteristics has two kinds of excitation regions, so we show visualized animation of the downstream cylinder oscillating at reduced velocity of 2.4. You can see an oscillation of the downstream cylinder induced by a symmetric vortex-shedding. And we show the other visualized pictures at different and high velocity of 3.42. At this velocity, it's certified that a downstream cylinder oscillates due to the alternate Karman vortex shedding.



(16) We show the response curve and Strouhal numbers of the downstream cylinder with narrow gap of 0.5, also and a visualized animation at the reduced velocity of 3.44. In this figure, a downstream cylinder oscillates with half the Strouhal frequency, locking to **alternate Karman vortex shedding**.



(17) **You will note that** the response curve of the downstream cylinder with middle gap distance of 1, has two excitation regions; that is, the reduced velocities are from 1.5 to 2 and from 3 to 5. You will find that the excitation of small amplitude at low reduced velocity of 1.71, is due to the movement-induced excitation, accompanied with symmetrical vortex-shedding. On the other hand, the animation at high velocity 3.56 shows that this excitation with a large amplitude is induced by alternate Karman vortex shedding like this.



(18) Finally, excitation regions of the downstream cylinder are mapped against gap distance and reduced velocity. We classify the excitation region as three regions, namely,

1R: Excitation of very narrow gap is induced by alternate Karman vortex shedding accompanied with unstable limit-cycle oscillation.

2R: Excitation due to symmetrical vortices is limited to the regions of middle gap-distance from 0.75 to 2.

3R: Downstream cylinder experiences a buffeting phenomenon in the case of wide gap over 2.

CONCLUSIONS

Flow-induced in-line oscillation of two tandem circular cylinders was experimentally studied by free-oscillation testing in a wind tunnel. We carried out two kinds of free-oscillation tests of two tandem cylinders. We measured the response amplitudes and the Strouhal frequency as functions of the gap distance. The flow around the cylinders was visualized by the smoke-wire method, also.

- (1) The in-line oscillation induced by symmetric vortex shedding of **the upstream cylinder** appears in all experimental ranges of gap. At wide gap-ratios $s = 2$ to 3 , the other excitation of $Vr = 3$ to 3.5 is induced by alternate Karman vortex shedding like a single cylinder.
- (2) The in-line excitation of **the downstream cylinder** is confirmed to be induced by alternate Karman vortex shedding for the narrow gap distances $s = 0.3$ to 0.75 , and the other excitation due to symmetric vortices are limited to gap distances of $s = 0.75$ to 2 . Furthermore, **the downstream cylinder** experiences a buffeting phenomenon by wake-fluctuation of the upstream cylinder, when the gap distance is greater than 2.5 .

(19) Flow-induced in-line oscillation of two tandem circular cylinders was experimentally studied by free-oscillation testing in a wind tunnel. We carried out two kinds of free-oscillation tests of two tandem cylinders.

- (1) The in-line oscillation induced by symmetrical vortex shedding of **the upstream cylinder** appears in all experimental range of a gap. At wide gap-ratios $s = 2$ to 3 , the other excitation of $Vr = 3$ to 3.5 is induced by alternate Karman vortex shedding like a single cylinder.
- (2) The in-line excitation of **the downstream cylinder** is confirmed to be induced by alternate Karman vortex shedding for the narrow gap distance $s = 0.3$ to 0.75 , and the other excitation due to symmetrical vortices are limited to gap distance ratios $s = 0.75$ to 2 .

Furthermore, the downstream cylinder experiences a buffeting phenomenon by wake-fluctuation of the upstream cylinder, when the gap distance is greater than $s = 2.5$.

**FLOW-INDUCED STREAMWISE OSCILLATION OF
TWO SQUARE CYLINDERS
IN TANDEM ARRANGEMENT**

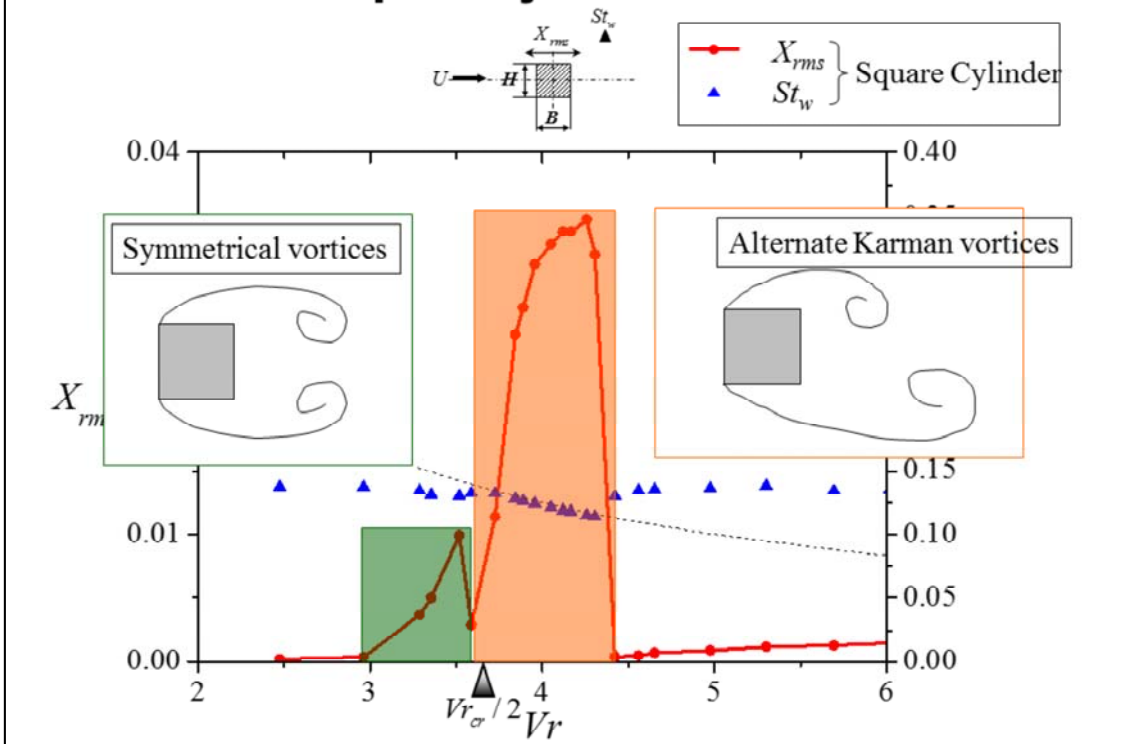
Atsushi OKAJIMA,
Professor, Kanazawa-Gakuin College
and
Takahiro KIWATA, Satoru YASUI, Yoshiki MORI, Shigeo KIMURA

***Keywords:* flow-induced oscillation, streamwise oscillation, two square cylinders, tandem arrangement, flow visualization**

(1) Thank you, Mr. Chairman.

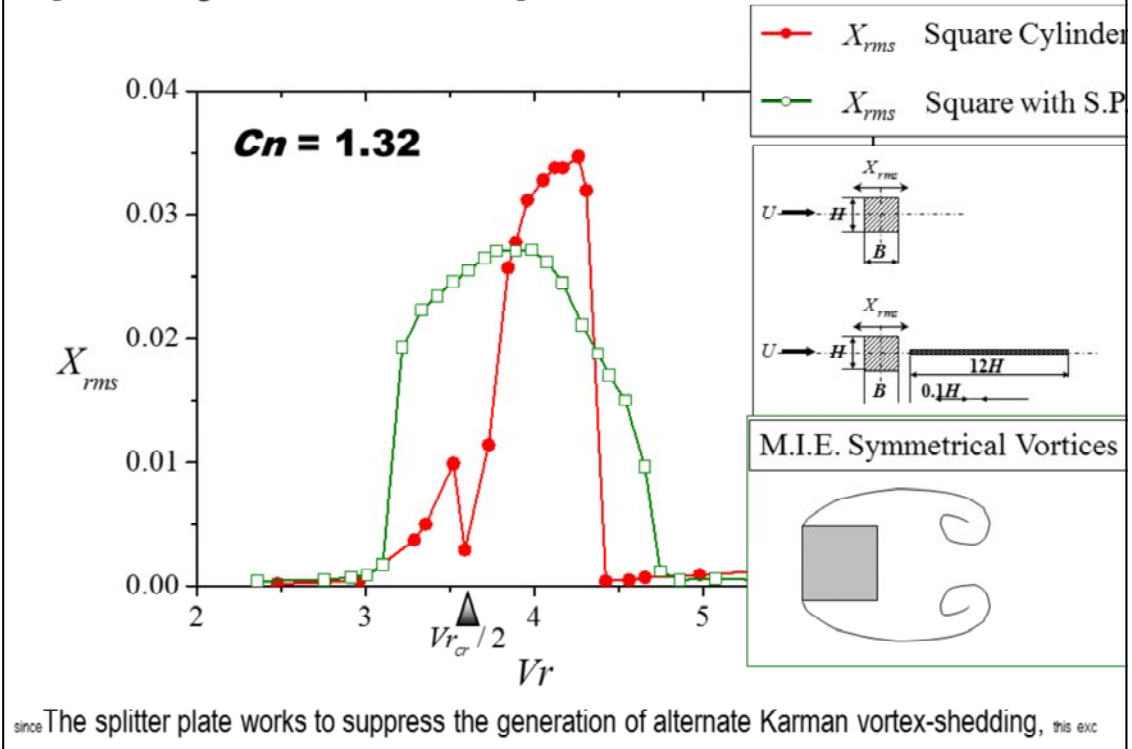
I would like to talk about flow-induced streamwise oscillation of two square cylinders in tandem arrangement.

PREVIOUS STUDIES : Flow-Induced Streamwise Oscillation of a Square Cylinder



(2) As fundamental information, I'll show you results of response amplitude for a single square cylinder obtained by the previous studies. The streamwise oscillations occur in two regions near half of the resonance velocity; the one at lower velocity is termed the first excitation region due to symmetrical vortices, and the other excitation at high velocities is termed the second excitation region due to alternate Karman vortices.

PREVIOUS STUDIES : Flow-Induced Oscillation of a Square Cylinder with a Splitter Plate



(3) When we insert a splitter plate in the wake of a square cylinder, the response amplitudes of the streamwise oscillation are plotted with the green line. Since the splitter plate works to suppress the generation of alternate Karman vortex-shedding, this excitation is apparently due to the movement-induced excitation accompanied by symmetrical vortices.

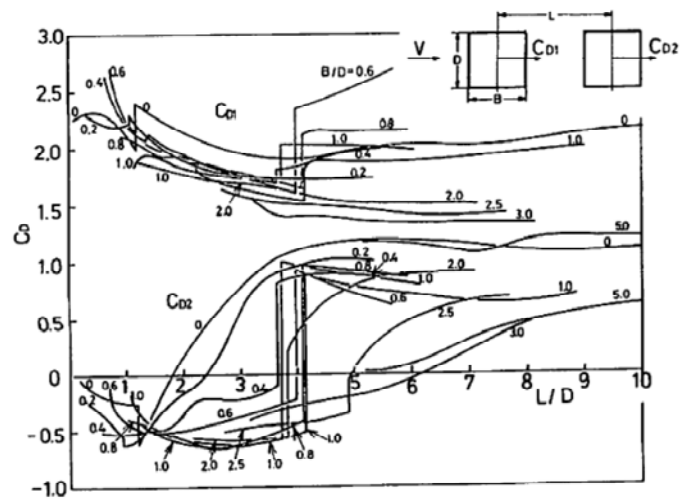
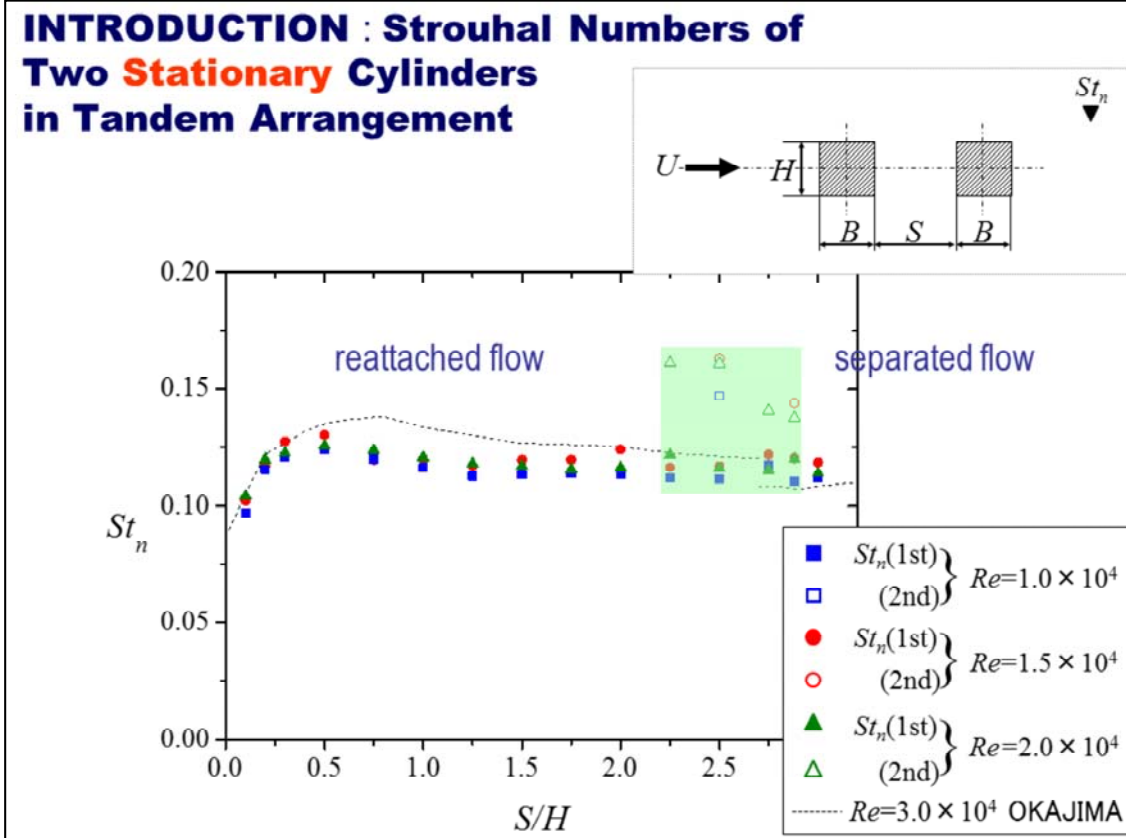


Figure 37. Drag coefficients of two tandem rectangular cylinders with various side ratios, B/D , of 0, 0.2, 0.4, 0.6, 0.8, 1.2, 2.5, 3, and 5 [28, 51, 52].



(4) We show the results of the Strouhal number behind two stationary cylinders in tandem arrangement against the gap-distance as the fundamental data of two tandem cylinders. You can see that the Strouhal number changes with spreading the gap-distance. The values of the Strouhal number seem to be close, slightly influenced by the Reynolds number. Only at these gap distances, the values of Strouhal number scatter, since the flow pattern changes from a reattached flow to a separated flow.

PARAMETERS FOR FLOW-INDUCED OSCILLATION

Reduced velocity

$$Vr = \frac{U}{f_c H}$$

Critical reduced velocity

$$Vr_{cr} = \frac{U}{f_n H}$$

Reduced mass-damping parameter

$$Cn = \frac{2m\delta}{\rho H^2 L}$$

Non-dimensional gap distance between two cylinders

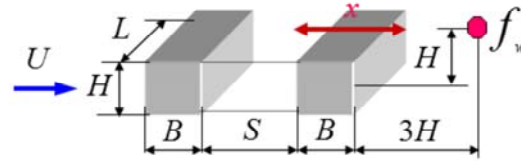
$$s = \frac{S}{H}$$

Non-dimensional response amplitude of a cylinder

$$X_{rms} = \frac{x_{rms}}{H}$$

Strouhal Number of Vortex-Shedding

$$St_w = \frac{f_w H}{U}$$



H : Cylinder height

f_c : Characteristic frequency of a cylinder

f_n : Natural vortex shedding frequency for a stationary cylinder

f_w : Frequency of a wake

L : Span length of a cylinder

m : Mass of a cylinder per unit span length

U : Uniform flow velocity

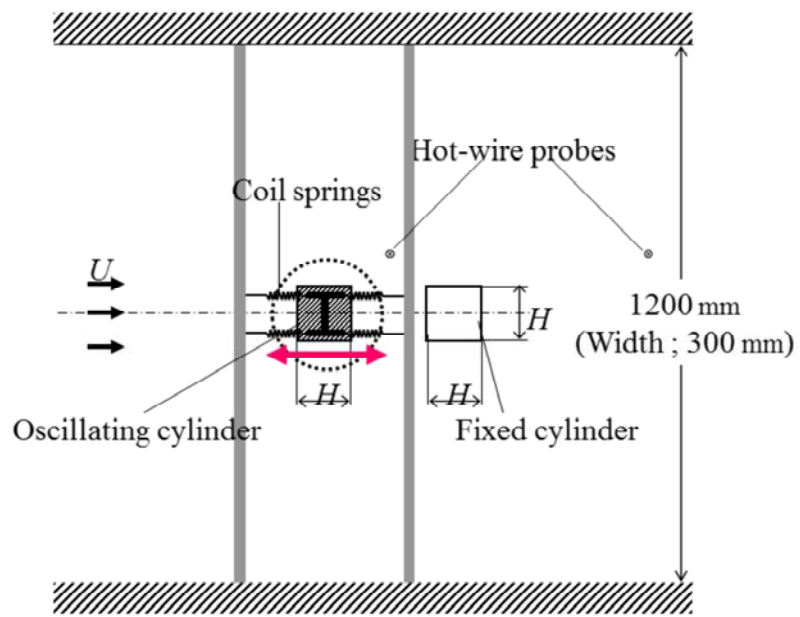
x_{rms} : Root-mean square response amplitude of a cylinder

δ : Logarithmic decrement of the structure damping parameter of a cylinder

ρ : Fluid mass density

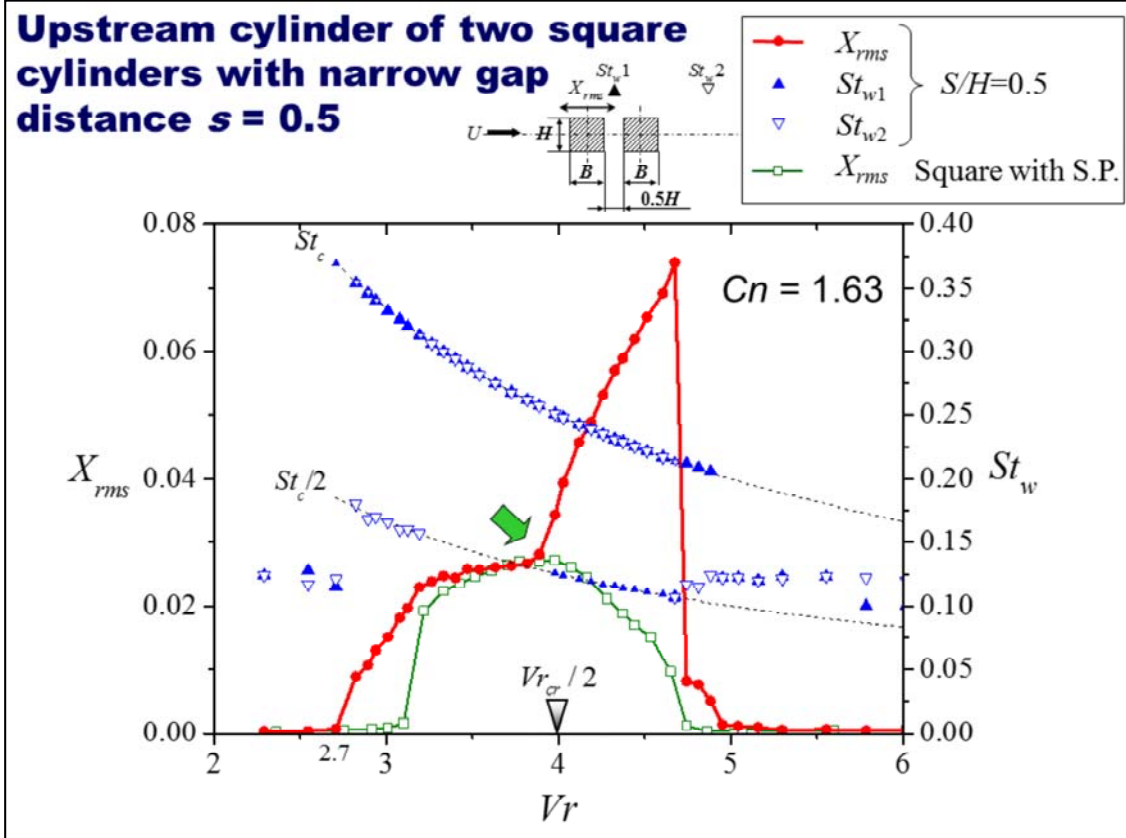
(5) When we carried out free-oscillation tests of two square cylinders in tandem arrangement, the reduced velocity, the critical reduced velocity, the reduced mass-damping parameter Cn and the gap distance S between the upstream and downstream cylinders, are important parameters. So we measured the non-dimensional response amplitudes and Strouhal number of the vortex-shedding as the experimental results.

Schematic View of a Test Section for Experiments

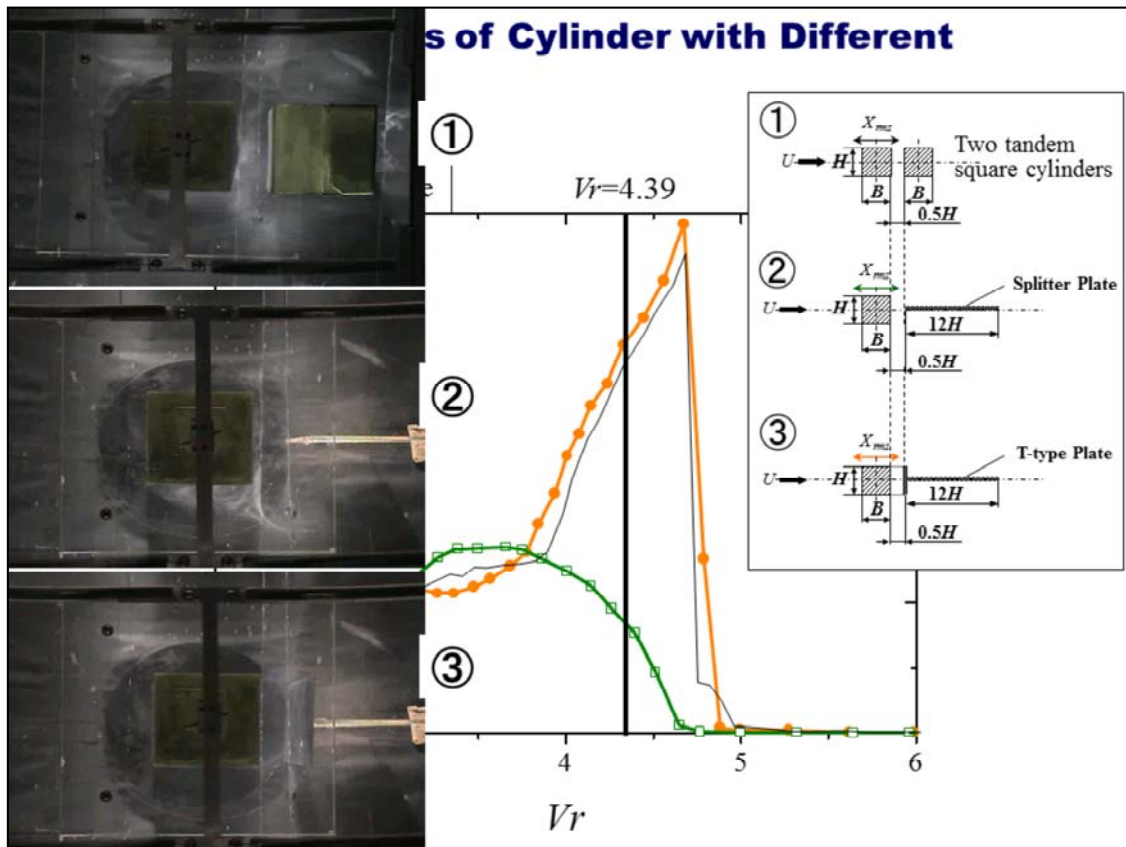


The first case: the upstream cylinder elastically supported to move in the streamwise direction, is free to oscillate, while the downstream cylinder is fixed.

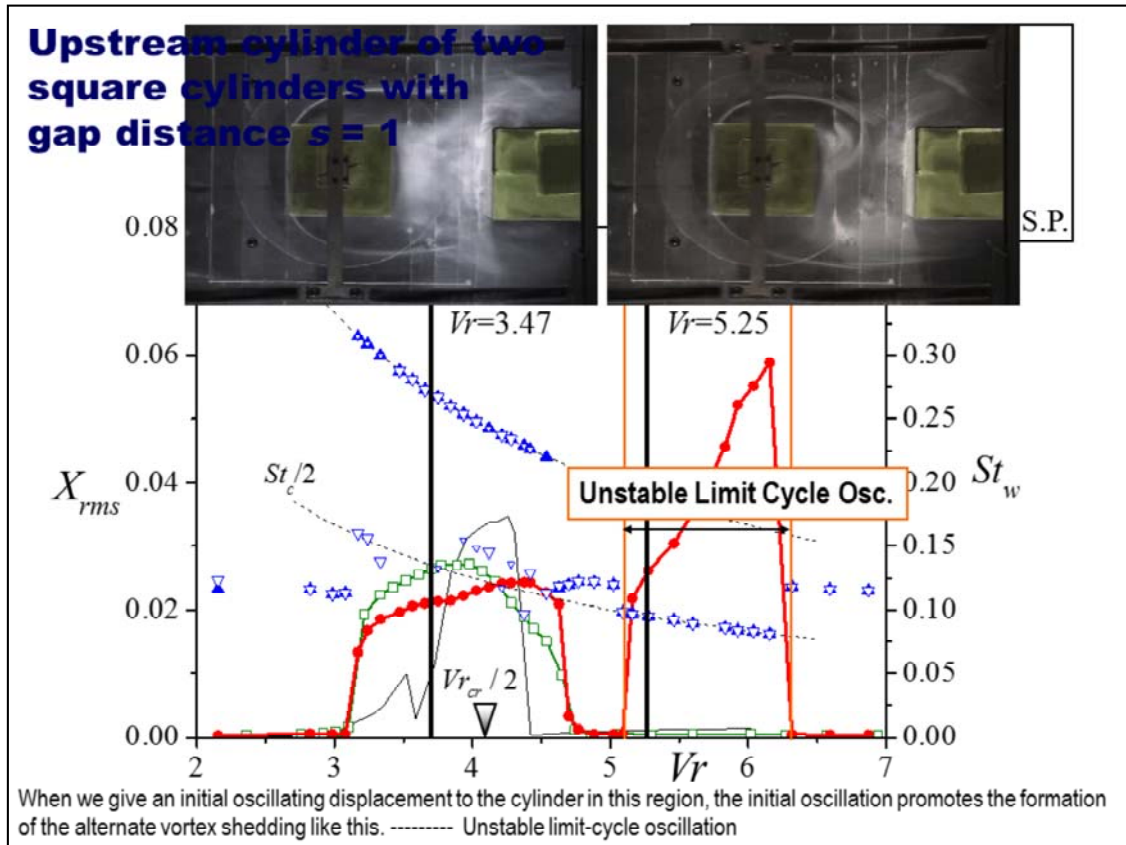
(6) Well I'll show a schematic view of a test section for experiments: In this case, the upstream cylinder elastically supported to move in the streamwise direction, is free to oscillate, while the downstream cylinder is fixed.



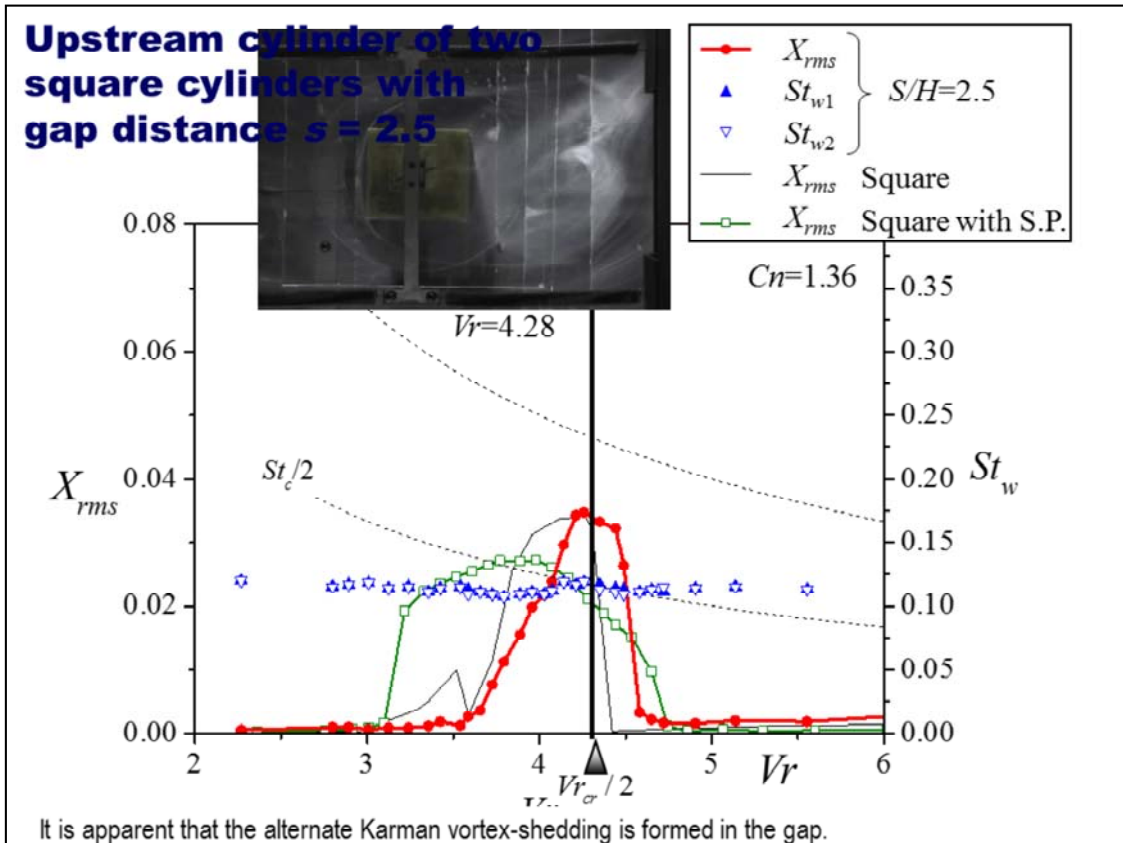
(7) I'll show an example of the response curve of the streamwise oscillation of the upstream cylinder for narrow gap distance 0.5. Streamwise oscillations begin to occur near $Vr = 2.7$, and the amplitude changes along the curve of the cylinder with a splitter plate with increasing the reduced velocity, and still the value of amplitude abruptly grows up to a maximum value of about 0.08. It is noted that this response curve is quite different from that of the cylinder with a splitter plate for this region. Why this excitation occurs, having very large amplitudes.



(8) Then we tried to carry out some **additional** experiments: (1) This response curve is for the case of two tandem cylinders. This curve is for the case of inserting a splitter plate, and when we employed a T-type splitter, this curve is very similar with that of two tandem cylinders. So, I'll show the visualized animations for each case: It is evident that the flows of case 1 and 3 seem to be almost similar, and you can see the wake approaches the side surfaces of the downstream cylinder or the edges of T-type splitter, **which produces a narrow wake**, while you can see the symmetrical vortices of the case 2 is comparatively wide. However, the excitation of all cases is due to movement-induced excitation accompanied by the symmetrical vortices.

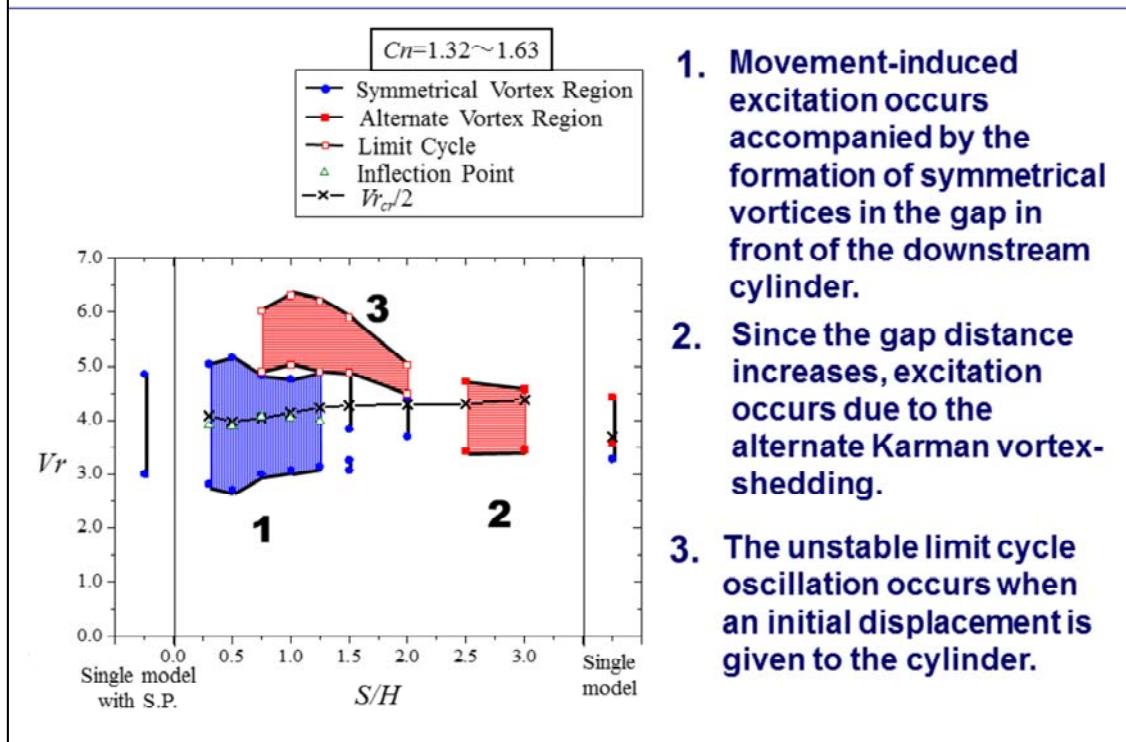


(9) This figure is the results of the upstream cylinder for a gap distance 1. There are separately two excitation regions. The response curve of the first excitation region is very similar to that of the cylinder with a splitter plate. So we show the visualized animations when the upstream cylinder oscillates in this region. A pair of symmetrical vortices form in the gap. And you can find the other excitation, that is called to be an unstable limit-cycle-oscillation. When we give an initial oscillating displacement to the cylinder in this region, the initial oscillation promotes the formation of the alternate vortex shedding, which results in establishing of the so-called unstable limit-cycle oscillation.



(10) This is the results of the streamwise oscillation of the upstream cylinder, when the gap distance becomes relatively wide, that is 2.5. Only one excitation region appears. The visualized animation shows that the upstream cylinder oscillates in the streamwise direction at $Vr = 4.28$, and it is apparent that the alternate Karman vortex-shedding is formed in the gap between the cylinders.

Map of Excitation Regions of Upstream Cylinder as Functions of Gap Distance s and Reduced Velocity Vr



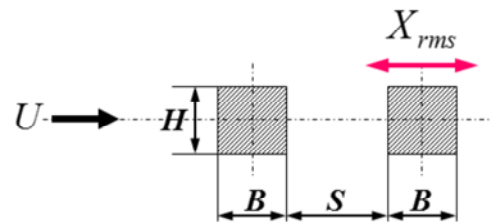
(11) Then we try to make a map of excitation regions of the upstream cylinder as functions of gap distance and reduced velocity. There are three excitation regions.

(1st Region) Movement-induced excitation **occurs** accompanied by the formation of symmetrical vortices in the gap in front of the downstream cylinder.

(2nd Region) When the gap distance increases, oscillation occurs due to the alternate vortex-shedding.

(Last Region) The unstable limit cycle oscillation occurs when an initial displacement is given to the cylinder.

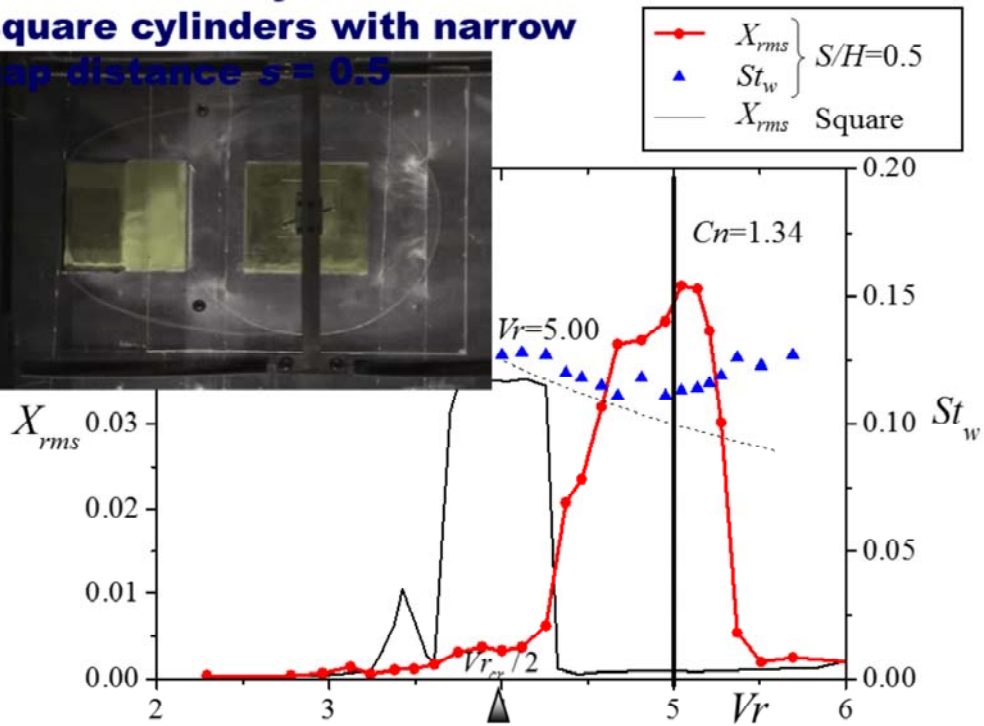
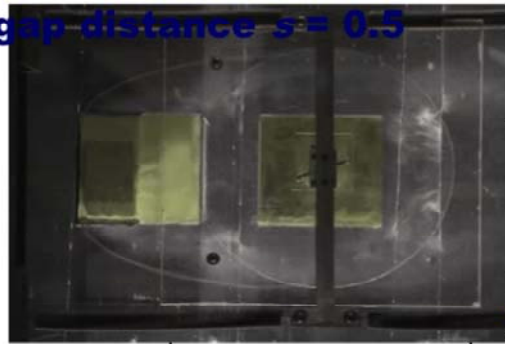
RESULTS : Flow-Induced Streamwise Oscillation of the Downstream Cylinder



The downstream cylinder is free for streamwise oscillation and the upstream cylinder is fixed.

(12) Next we present the results when the downstream cylinder is free for streamwise oscillation and the upstream cylinder is fixed.

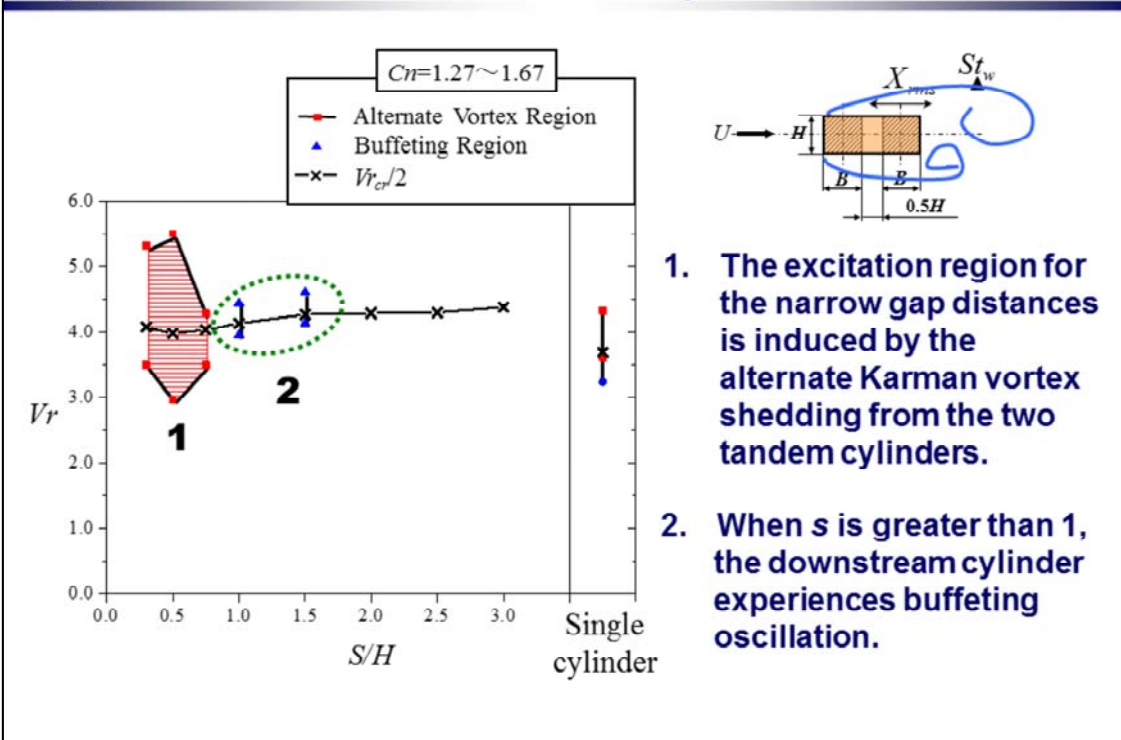
Downstream cylinder of two square cylinders with narrow gap distance $s = 0.5$



The downstream cylinder oscillates due to the alternate Karman vortex-shedding from two cylinders.

(13) We show the response curve of the downstream cylinder for narrow gap distance, 0.5 . You can find only one excitation region is quite different from those of a single cylinder; Then, we show a visualized picture for the downstream cylinder oscillating in the streamwise direction. The downstream cylinder oscillates due to the alternate Karman vortex-shedding from two cylinders.

Map of excitation regions of downstream cylinder as functions of gap distance s and reduced velocity V_r



(14) Finally, we show a map of excitation regions of the downstream cylinder against gap distance and reduced velocity.

(1) The excitation region for the narrow gap distances is induced by the alternate Karman vortex shedding from the two tandem cylinders, connected by the dead water region like this.

(2) When s is greater than 1, the downstream cylinder experiences buffeting oscillation.

CONCLUSIONS

- 1. When the upstream cylinder is free to oscillate, there are two excitation regions: the first region is due to the movement-induced excitation accompanied by the symmetrical vortex shedding, while the second region is due to the vortex excitation by the alternate Karman vortex shedding accompanied with the unstable limit-cycle oscillation.**
- 2. For the range of high velocity, the initial displacement of the upstream cylinder promotes the periodic reattachment of the separated shear layers of the upstream cylinder to the downstream one, which results in establishing of the unstable limit-cycle oscillation.**
- 3. For wide gaps over 2.5, the excitation region of the upstream cylinder occurs, being due to the alternate Karman vortex shedding, and it resembles the streamwise oscillation of a single cylinder, since the influence of the downstream cylinder becomes small.**
- 4. When the downstream cylinder is free to oscillate for the narrow gap distance, this excitation is apparently induced by the alternate Karman vortex-shedding from the two cylinders**
- 5. When the gap-distance is greater than 1, the downstream cylinder experiences buffeting by wake fluctuation of the upstream cylinder.**

(15) Flow-induced in-line oscillation of two tandem circular cylinders was experimentally studied by free-oscillation testing in a wind tunnel. We carried out two kinds of free-oscillation tests of two tandem cylinders.

1. When the upstream cylinder is free to oscillate, there are two excitation regions: the first region is due to the movement-induced excitation accompanied by the symmetrical vortex shedding, while the second is due to the vortex excitation by the alternate Karman vortex shedding accompanied with the unstable limit-cycle oscillation.
2. For the range of high velocity, the initial displacement of the up-stream cylinder promotes the periodic reattachment of the separated shear layers of the upstream cylinder to the downstream one, which results in the unstable limit-cycle oscillation with large amplitude.
3. For wide gaps over 2.5, the excitation region of the upstream cylinder occurs, being due to the alternate Karman vortex shedding, and it resembles the streamwise oscillation of a single cylinder, since the influence of the downstream cylinder becomes small.
4. When the downstream cylinder is free to oscillate for the narrow gap distance, this excitation is apparently induced by the alternate Karman vortex-shedding from the two cylinders connected by the dead water region between them.
5. When the gap-distance is greater than 1, the downstream

cylinder experiences buffeting by wake fluctuation of the upstream cylinder.

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